

# An Intro to Ion Traps and the Berkeley ElectroStatic Trap



# Outline

- Short intro to ion traps.
- A novel ElectroStatic trap.
- Some possible applications.
- Project status.



# Why Ion Traps?

- Localization.
- Low Energy.
- Dilute System.

Can be used for...

- Quantum Information Measurements.
- Mass Measurements.
- Beyond SM high precision studies (EDM,  $\beta$  decay, APV,....).



# General Principles

- 3D trapping requires a potential minimum. Conveniently taken to be harmonic.

$$F \propto -\vec{r}$$

$$F = -\nabla U$$

$$\Phi = \frac{\Phi_0}{d^2} (Ax^2 + Bx^2 + Cz^2)$$

- In general  $\Phi$  can be a time-dependent function.
- Also note that this means no electrostatic traps.

$$\nabla^2 \Phi = 0 \Rightarrow A + B + C = 0$$



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Only saddles in 3D.  
“Ernshaw’s Theorem”  
(more on this later)

$$A + B + C = 0$$



# General Principles

$$\nabla^2 \Phi = 0 \Rightarrow A + B + C = 0$$

Take rotational symmetry to get  $A=B=1$ ,  $C=-2$ .

$$\Phi = \frac{\phi_0}{d^2} (x^2 + y^2 - 2z^2) = \frac{\phi_0}{d^2} (\rho^2 - 2z^2)$$

Hyperbolic electrode surfaces:

$$\rho^2 - 2z^2 = r_0^2$$

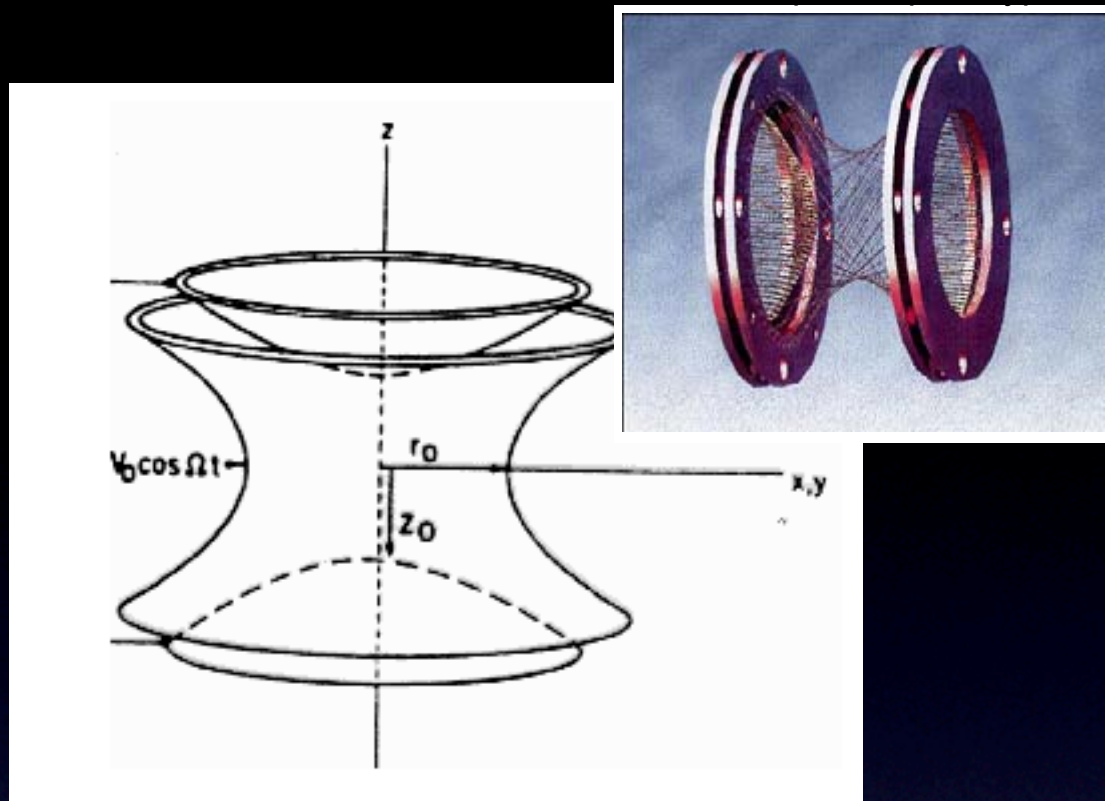
$$\rho^2 - 2z^2 = -2z_0^2$$

$$d^2 = r_0^2 + 2z_0^2$$

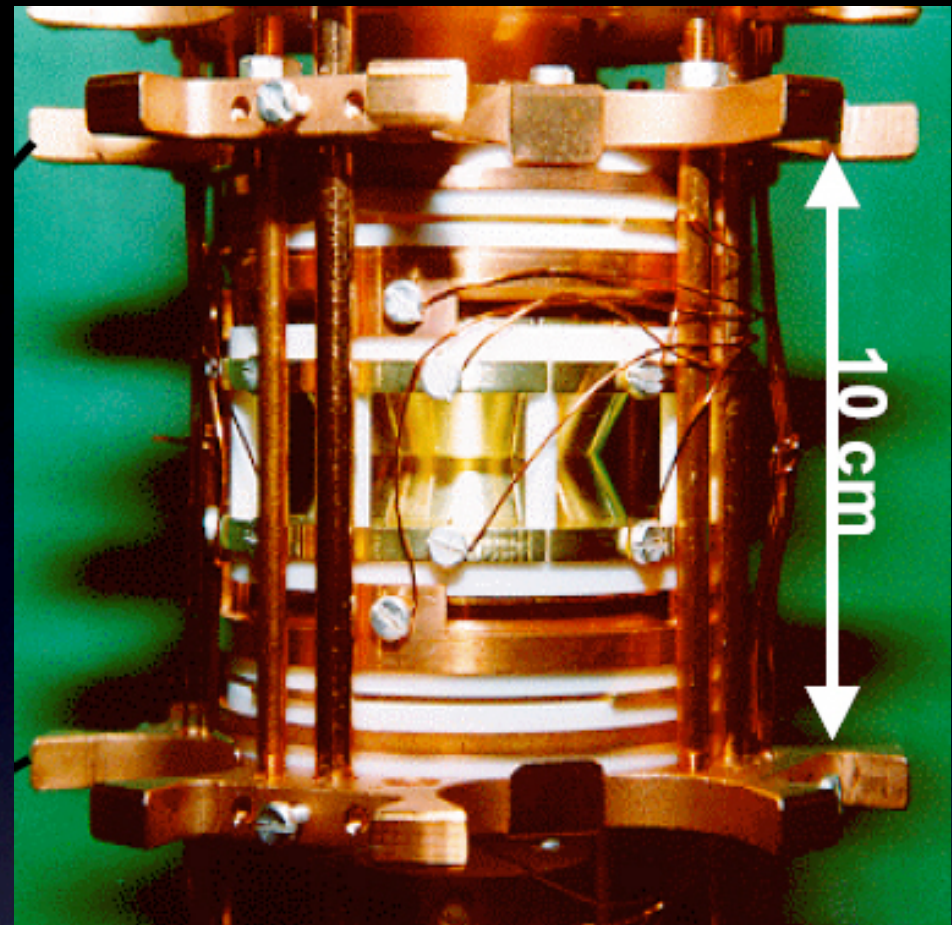
Two general types of traps:

- Penning trap: superimpose a magnetic field along the z axis.
- Paul trap: use time varying (RF) electric fields.
- Linear variant of the Paul trap.

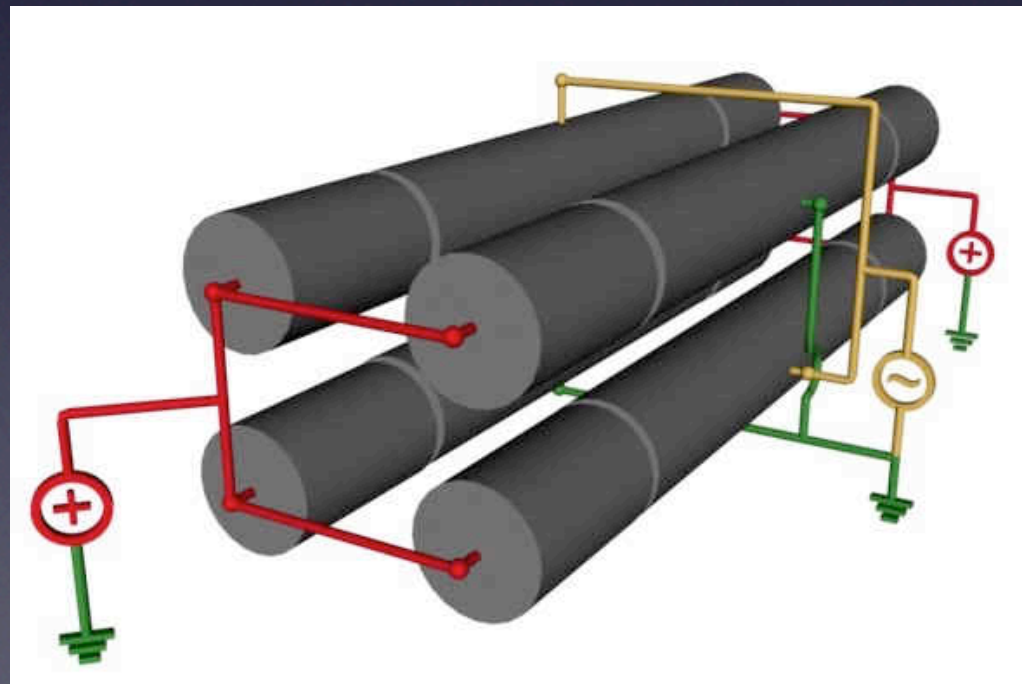




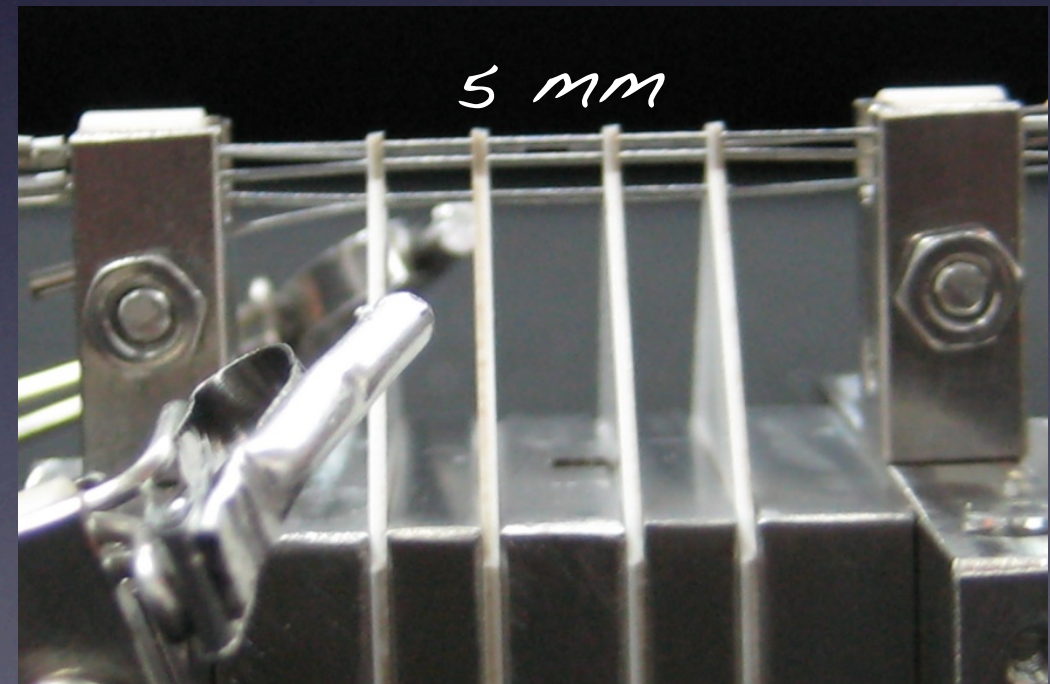
RF Paul Trap



Penning trap at ISOLDE



Linear Paul Trap



Linear Paul Trap at WI



# Some Math (sorry...)

## Paul Traps

$$\Phi = \frac{U_0 + V_0 \cos \Omega t}{2d^2} (r^2 - 2z^2)$$

Applied between ring  
and endcaps

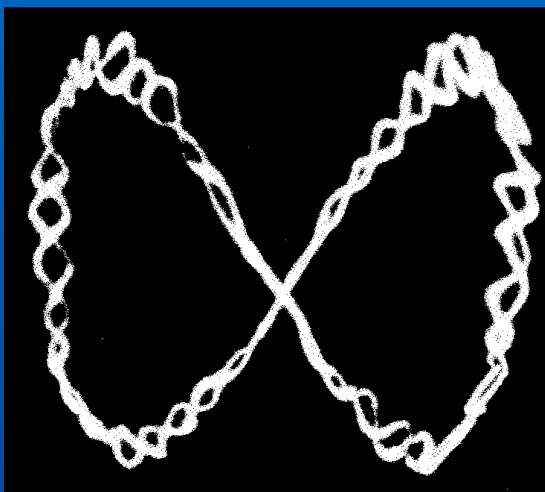
EOM (homogeneous Mathieu):

$$\frac{\partial^2 u}{\partial t^2} = \frac{Q}{Md^2} (U_0 + V_0 \cos \Omega t) u$$

$$\frac{\partial^2 u}{\partial t^2} + (a - 2q\tau)u = 0$$

$$(\tau = \Omega t, a_x = a_y = -2a_z = -4QU_0/Md^2\Omega^2,$$

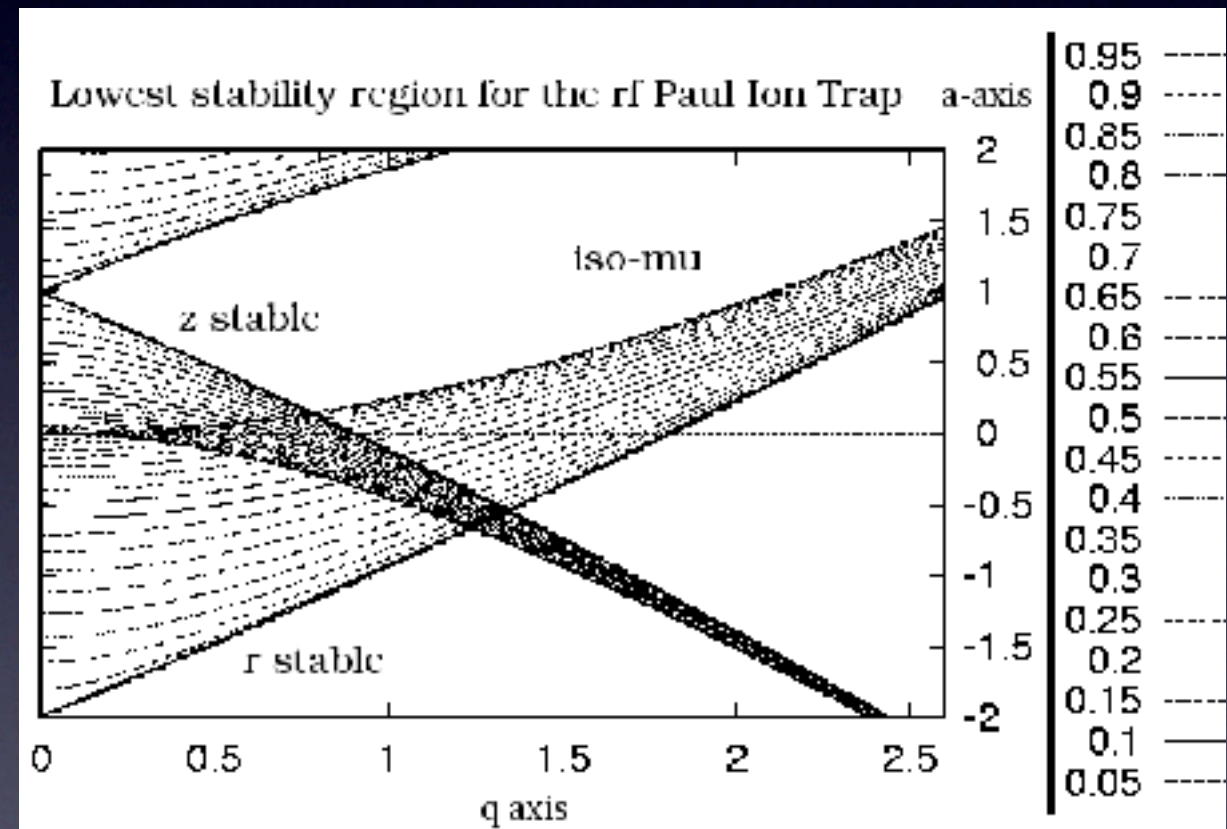
$$q_z = q_y = 2q_x = 2QV_0/Md^2\Omega^2)$$



Slow *macromotion*  
super imposed with  
RF drive frequency

Wuerker et al., J. Appl Phys 30, 342 (1959)

## Stability Diagram



$$u_i(t) \simeq A \left( 1 - \frac{q_i}{2} \cos \Omega t \right) \cos \omega_i t$$

$$\omega_i = \beta_i \Omega / 2, \beta_i \simeq a_i + q_i^2 / 2$$



# Some Math (sorry...)

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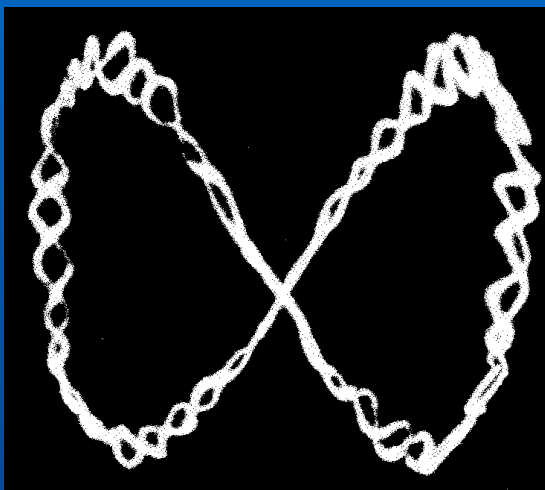
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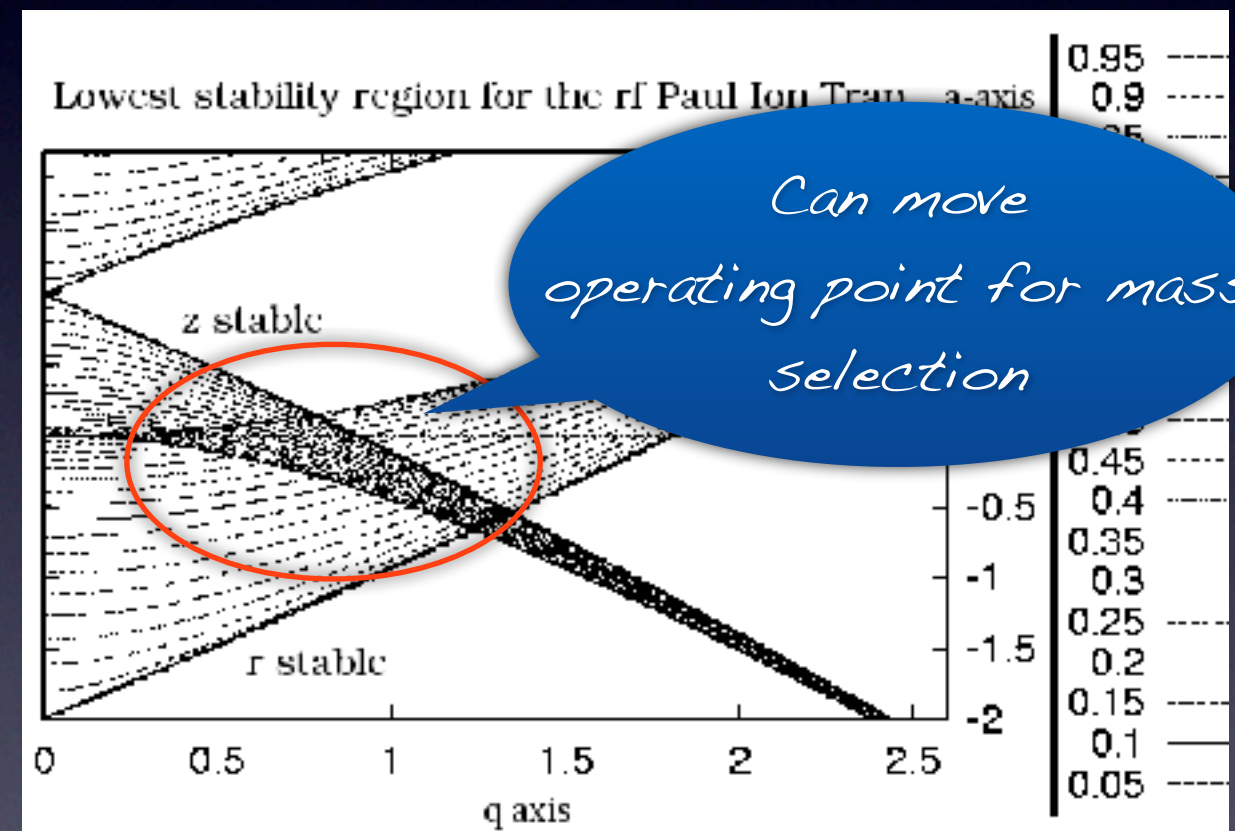
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$$\omega_i = \beta_i \Omega / 2, \beta_i \simeq a_i + q_i^2 / 2$$



# Some Math (sorry...)

## *Penning Traps*

$$\Phi = \frac{\Phi_0}{2d^2}(r^2 - 2z^2)$$

$$\vec{B} = (0, 0, B)$$

$$\vec{F} = -Q\Phi + Q(\vec{v} \times \vec{B})$$

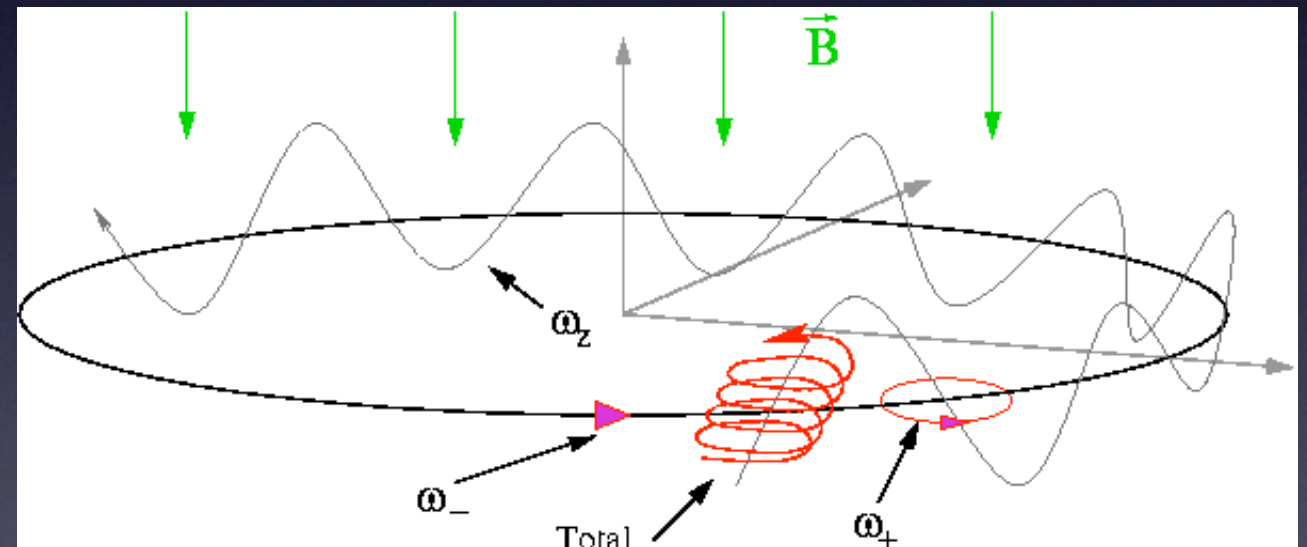
Apply a magnetic field +  
DC electric field

$$\frac{d^2x}{dt^2} - w_c^2 \frac{dy}{dt} - \frac{1}{2}w_z^2 x = 0$$

$$\frac{d^2y}{dt^2} + w_c^2 \frac{dx}{dt} - \frac{1}{2}w_z^2 x = 0$$

$$\frac{d^2z}{dt^2} + w_z^2 z = 0$$

$$w_z = \sqrt{\frac{2Q\Phi_0}{Md^2}}, \quad w_c = \frac{|QB_0|}{M}$$



$$w_+ = \frac{1}{2}(w_c + \sqrt{w_c^2 - 2w_z^2})$$
$$w_- = \frac{1}{2}(w_c - \sqrt{w_c^2 - 2w_z^2})$$



# Now What?

## *Detection*

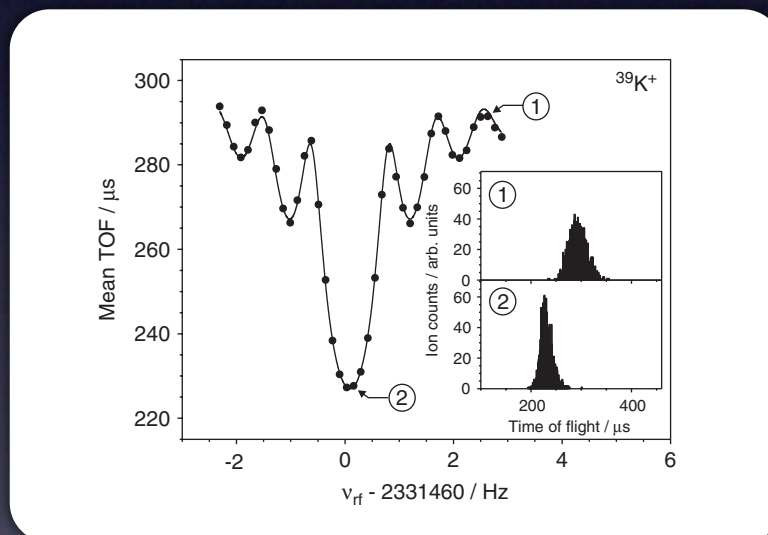
- Destructive:
  - Lower endcap voltage.
  - HV kick.
  - TOF-ICR (Penning Trap).



# Now What?

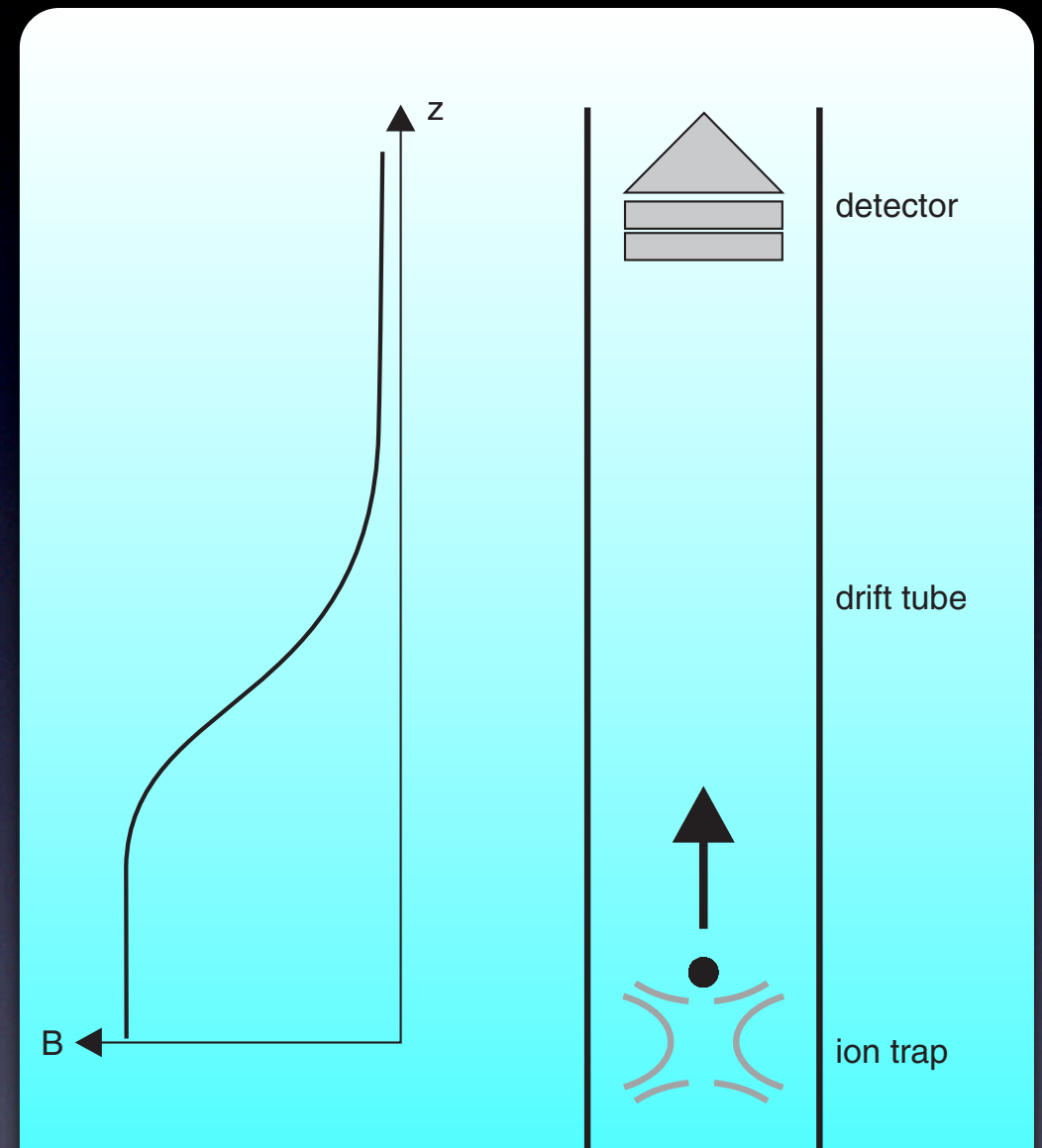
## Detection

- Destructive:
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  - HV kick.
  - TOF-ICR (Penning Trap).



Resolution of  $\sim 10^{-9}$  achieved for stable ions.

*K. Blaum, Phys. Rep. 425, 1 (2006)*



$$\vec{F} = -\vec{\mu}(\nabla \vec{B}) = -\frac{E_r}{B} \frac{\partial B}{\partial z} \hat{z}$$

$$T_{\text{tof}} = \int_0^{z_1} \sqrt{\frac{M}{2(E_0 - qU(z) - \mu B(z))}}$$



# Now What?

## *Detection*

- Destructive:
  - Lower endcap voltage.
  - HV kick.
  - TOF-ICR (Penning Trap).
- Non-Destructive:

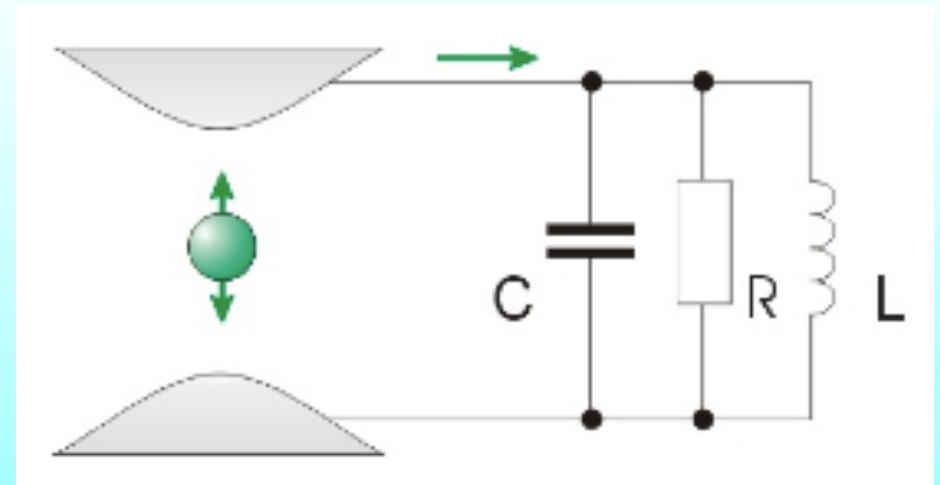


# Now What?

## *Detection*

- Destructive:
  - Lower endcap voltage.
  - HV kick.
  - TOF-ICR (Penning Trap).
- Non-Destructive:
  - Bolometric detection.

Moving ion = Moving  
charge = current



$$I = \Gamma \frac{q\dot{z}}{2z_0}$$

$$U_{noise} = \sqrt{4k_B T R \Delta\nu}$$

Also cools ions  
(resistive cooling)



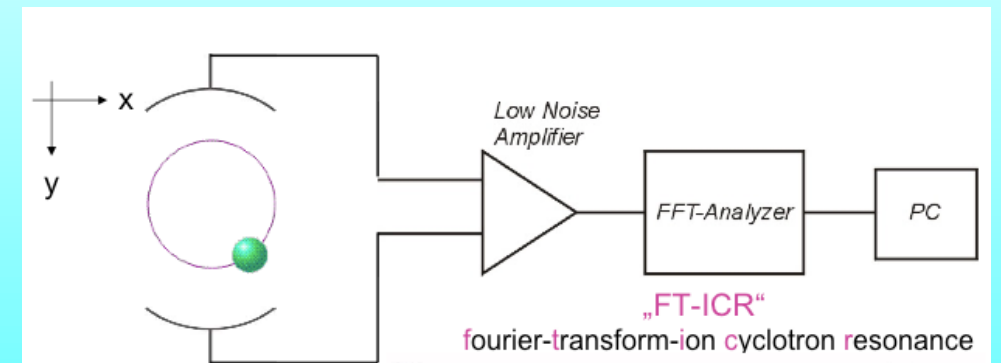
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## *Detection*

- Destructive:
  - Lower endcap voltage.
  - HV kick.
  - TOF-ICR (Penning Trap).
- Non-Destructive:
  - Bolometric detection.
  - Fourier Transform.
    - Broadband.
    - Resonant.

Moving ion = Moving  
charge = current

Detection of image  
charges on electrodes



More on this later

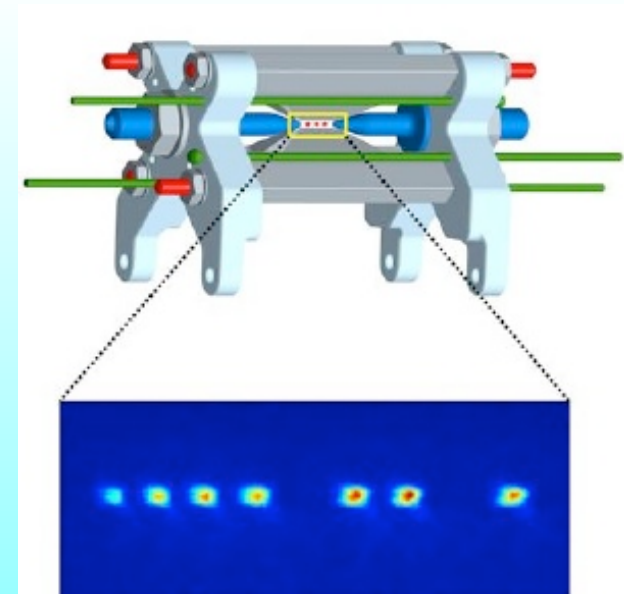


# Now What?

## *Detection*

- Destructive:
  - Lower endcap voltage.
  - HV kick.
  - TOF-ICR (Penning Trap).
- Non-Destructive:
  - Bolometric detection.
  - Fourier Transform.
    - Broadband.
    - Resonant.
  - Optical.

Moving ion = Moving  
charge = current



Easy, intuitive, coool, but:

- Expensive.
- Limited (mostly alkali-like ions).

Also, lets you do laser cooling.



# An Electrostatic Ion Beam Trap (EIBT) and (*some*) Possible Applications



# Ernshaw's Theorem

A collection of point charges cannot be maintained in a stable stationary equilibrium configuration solely by the electrostatic interaction of the charges.

Restatement of Gauss' Law (for free space)

$$\nabla \cdot E \propto \nabla \cdot F = -\nabla^2 \phi = 0$$

No local minima or maxima in free space (only saddle points).

Naively speaking → **No electrostatic ion traps**



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**Exceptions for:**

**Time varying & Magnetic fields.**

**Electronic correction.**

**Diamagnetic materials (aka “Floating Frog”).**



# Ernshaw's Theorem

A collection of point charges cannot be maintained in a stable stationary equilibrium configuration solely by the electrostatic interaction of the charges.

Rest

$$\nabla \cdot \mathbf{E}$$

(e)

$$= 0$$

No local minima of

(points).

Naively speaking -

**Exceptions for:**

**Time varying & Magnetic fields.**

**Electronic correction.**

**Diamagnetic materials (aka "Floating Frog").**





# An Electrostatic **Ion Beam** Trap (EIBT) and (*some*) Possible Applications

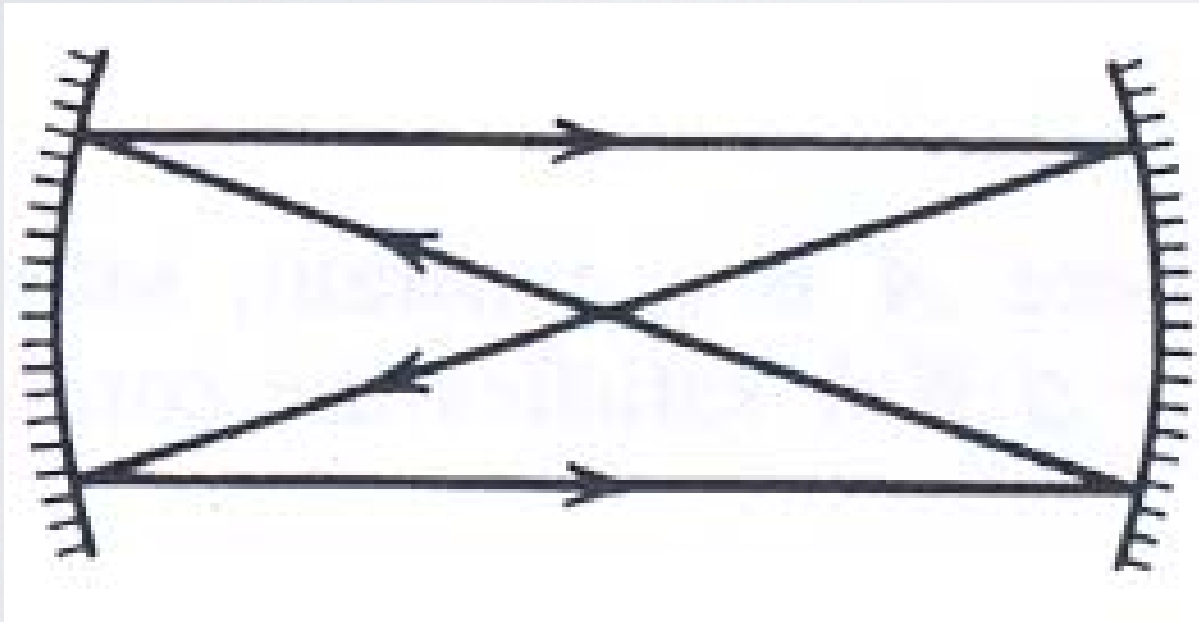


# But what about moving ions...

Ernshaw's theorem talks about stationary charges.

Moving charges in an electrostatic field actually “see” changing fields.

Trap design very similar to a resonant cavity for laser light.



*Optical Stability Condition*

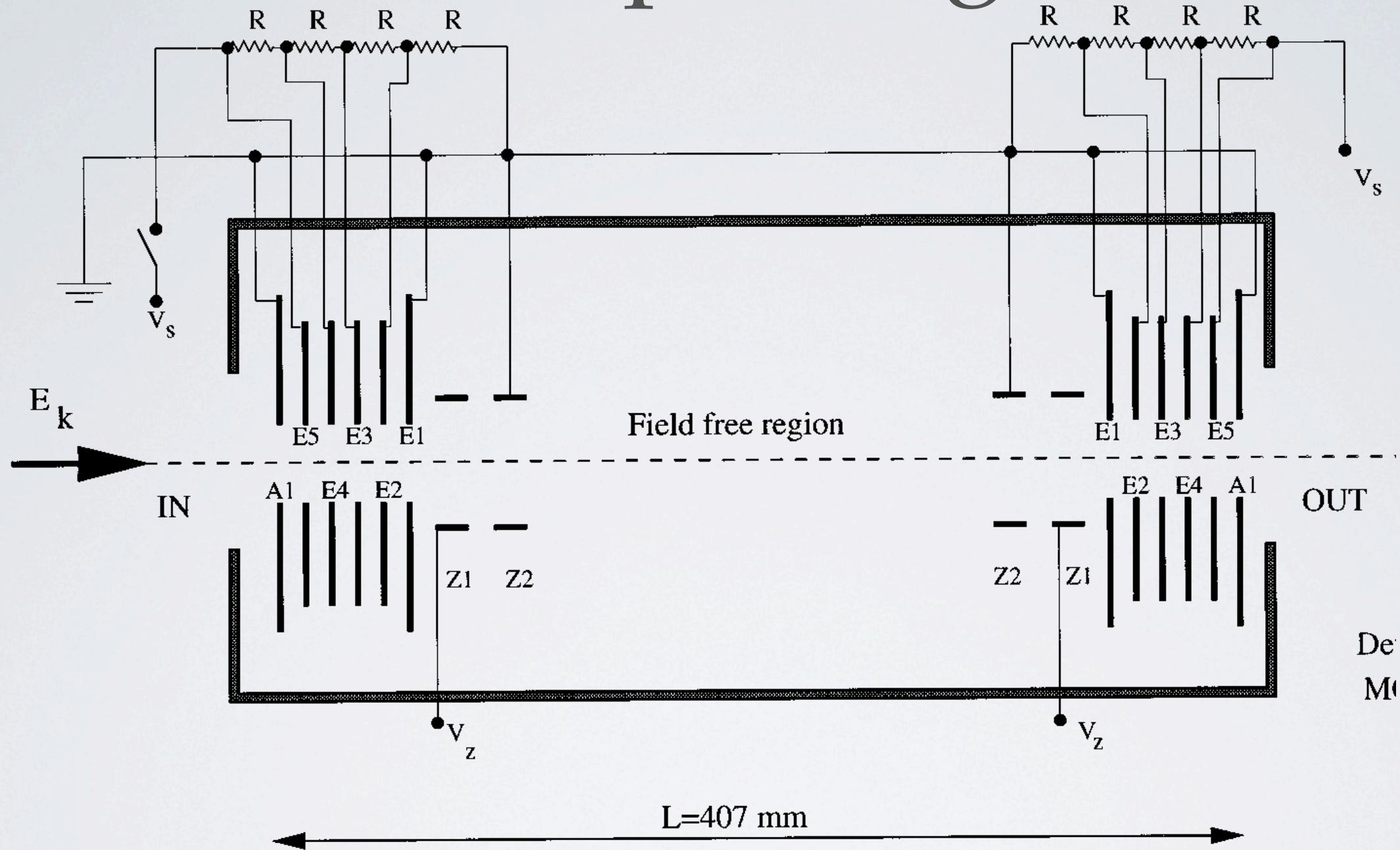
$$\frac{L}{4} \leq f \leq \infty$$

*Confining Potential Maximum*

$$q e V_s > E_k$$

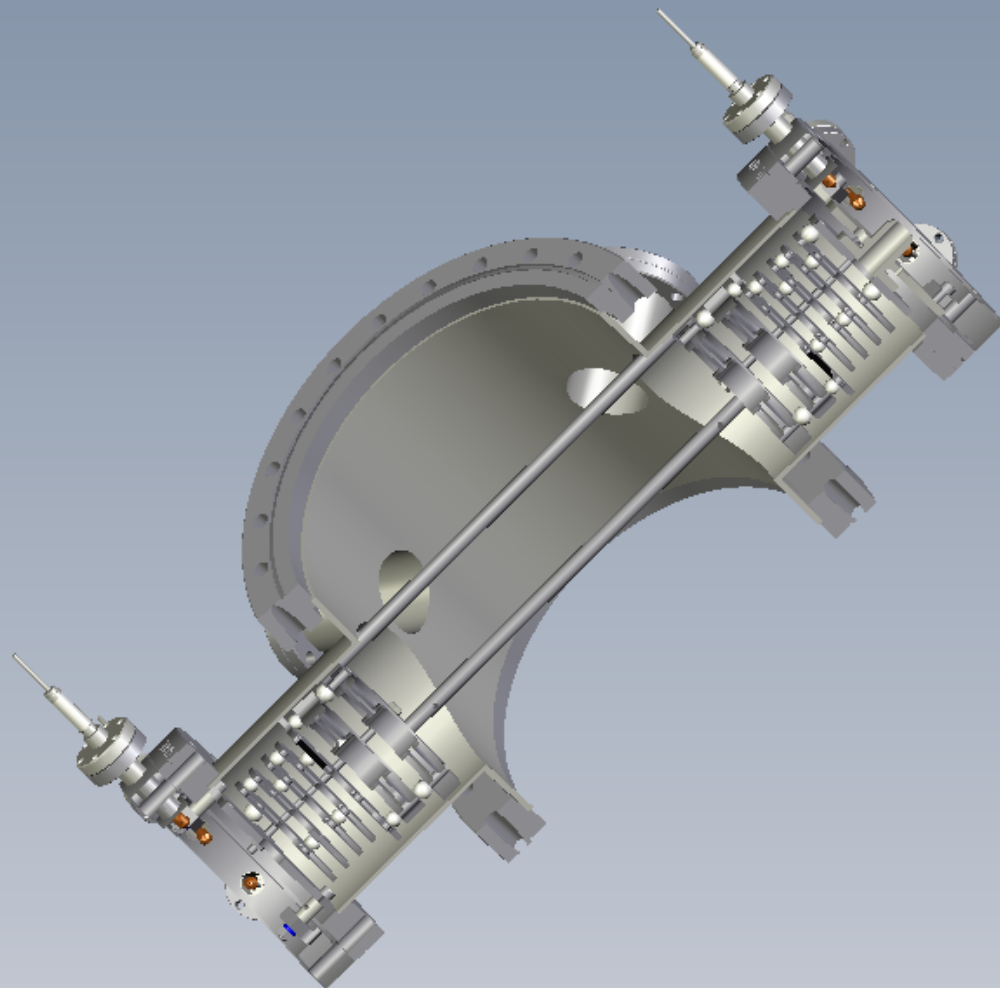
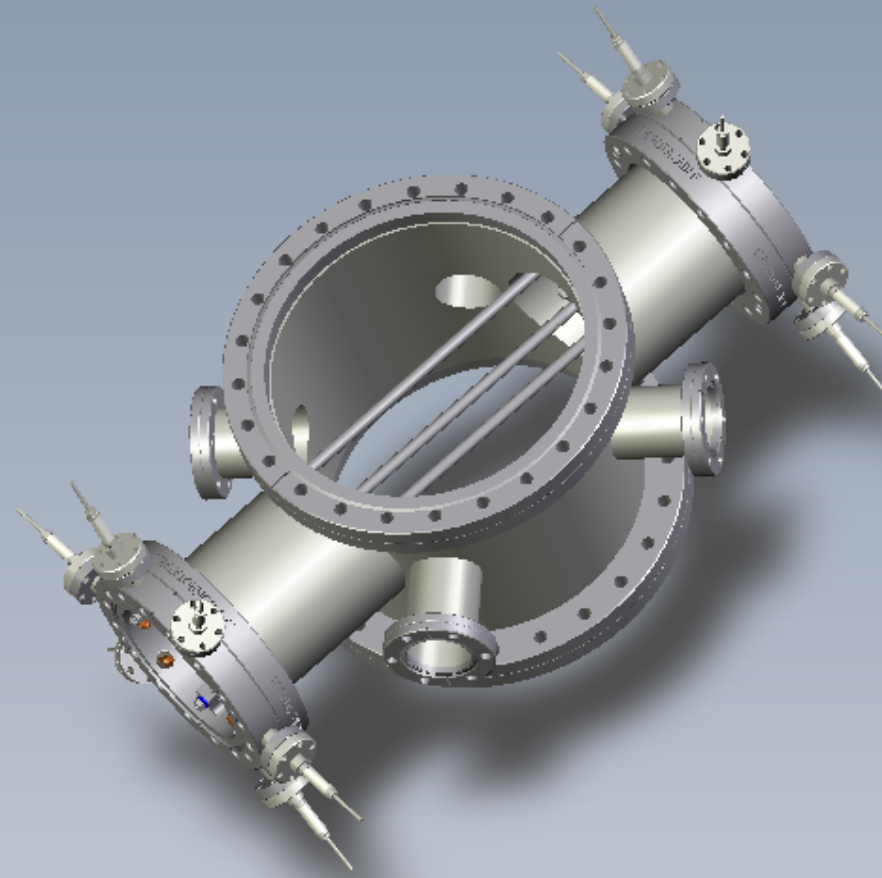


# Trap Design

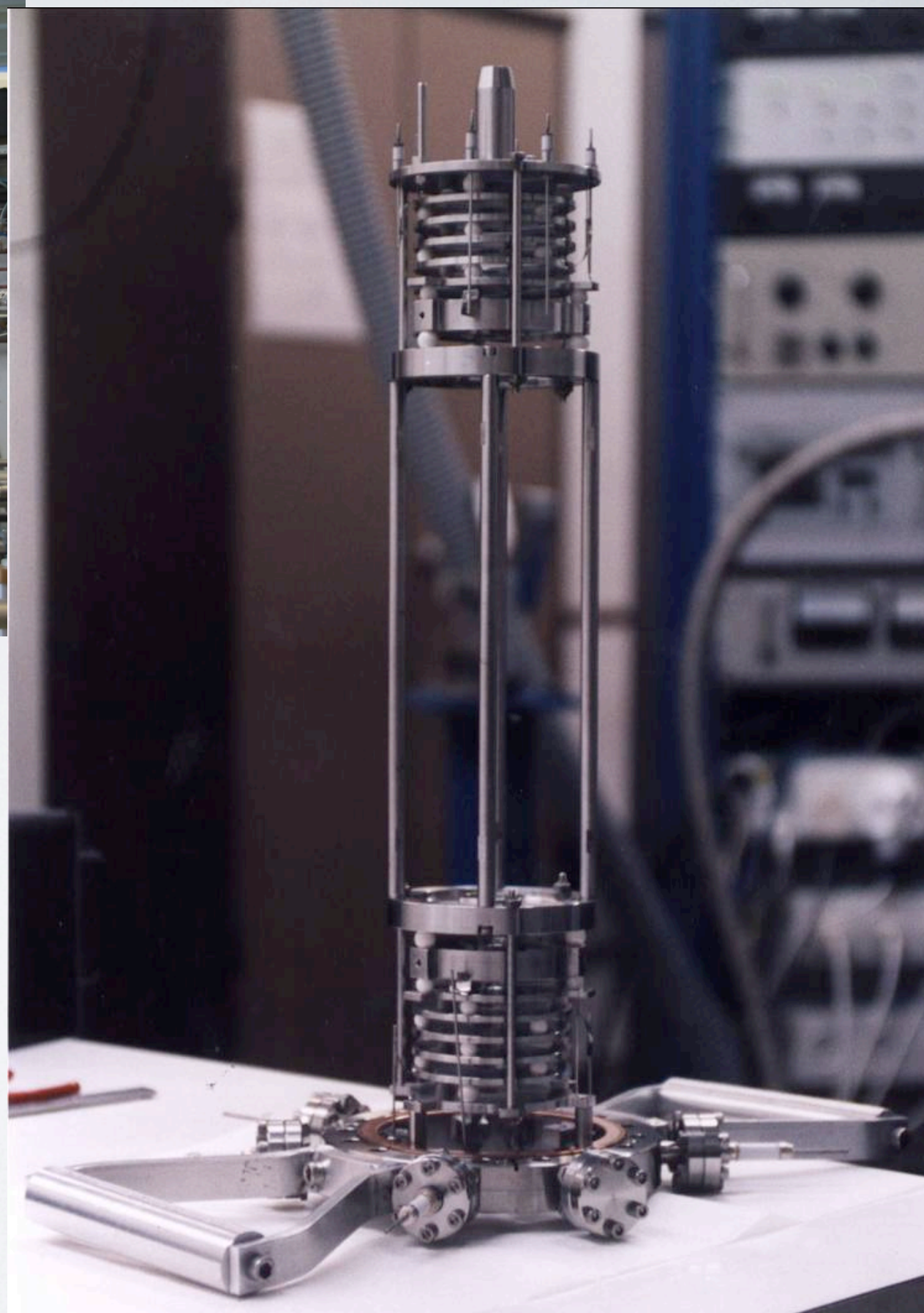
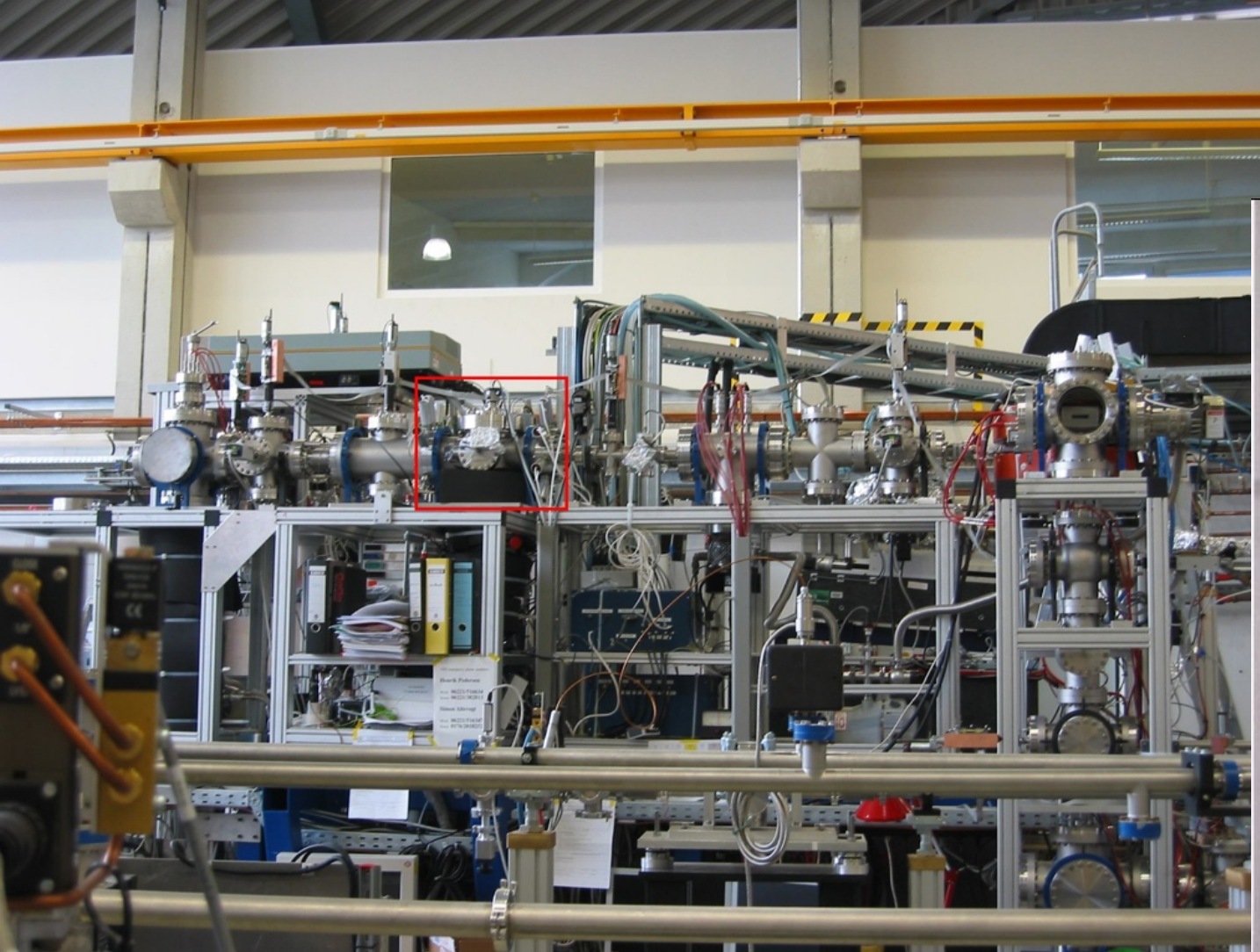


Realized at the Weizmann Institute (Israel). Also used in a cryogenic setup in Heidelberg.

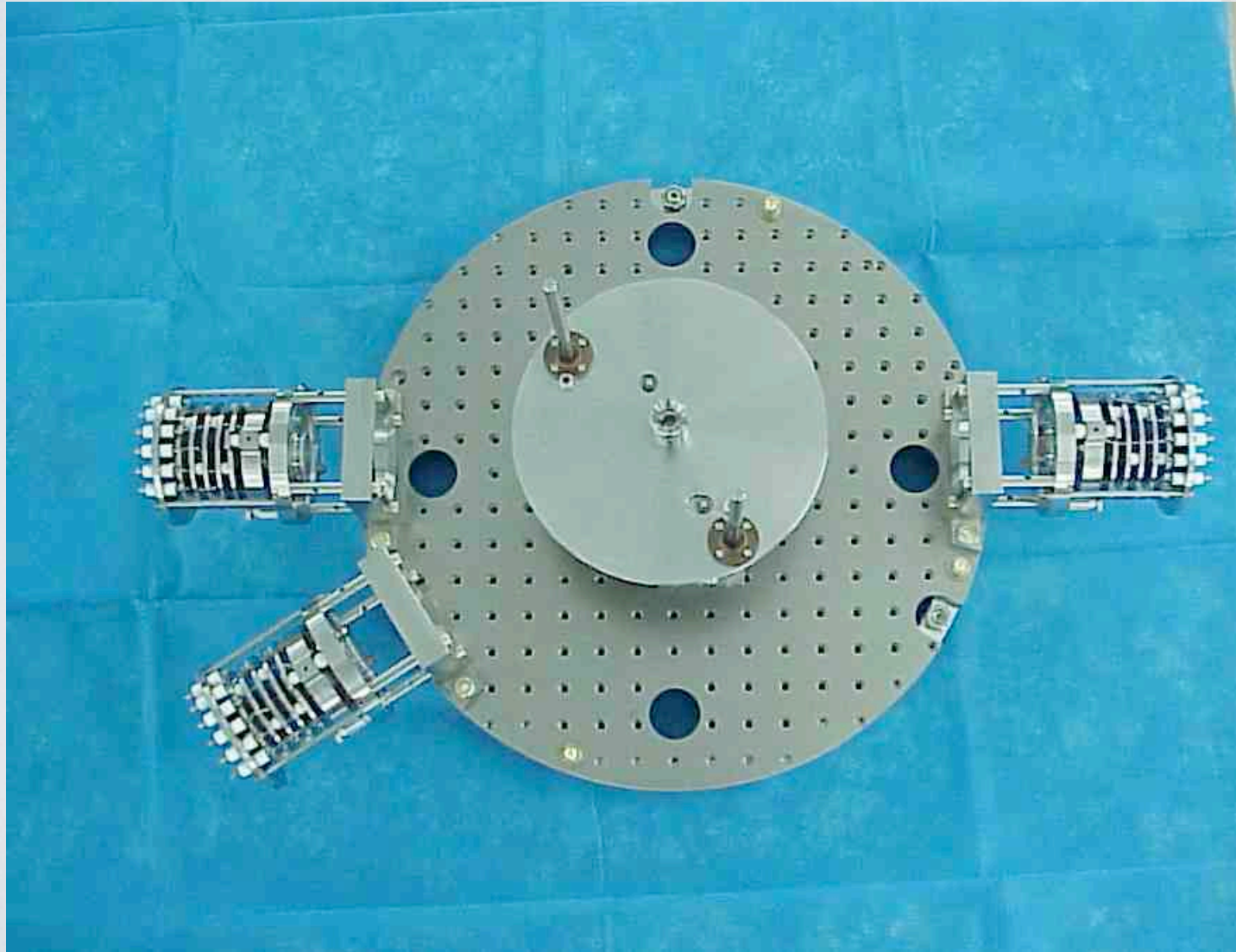




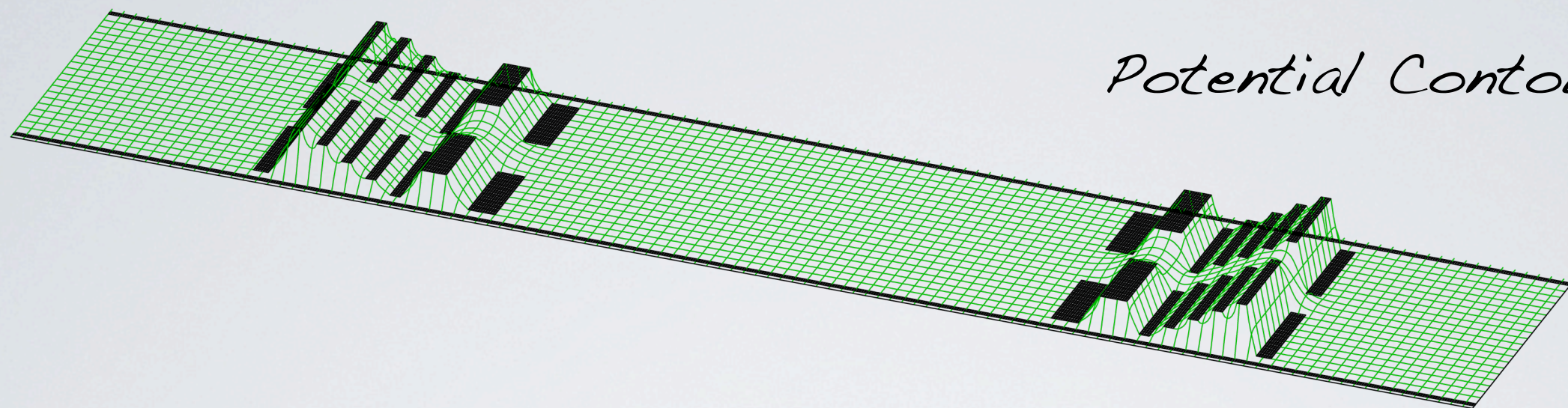




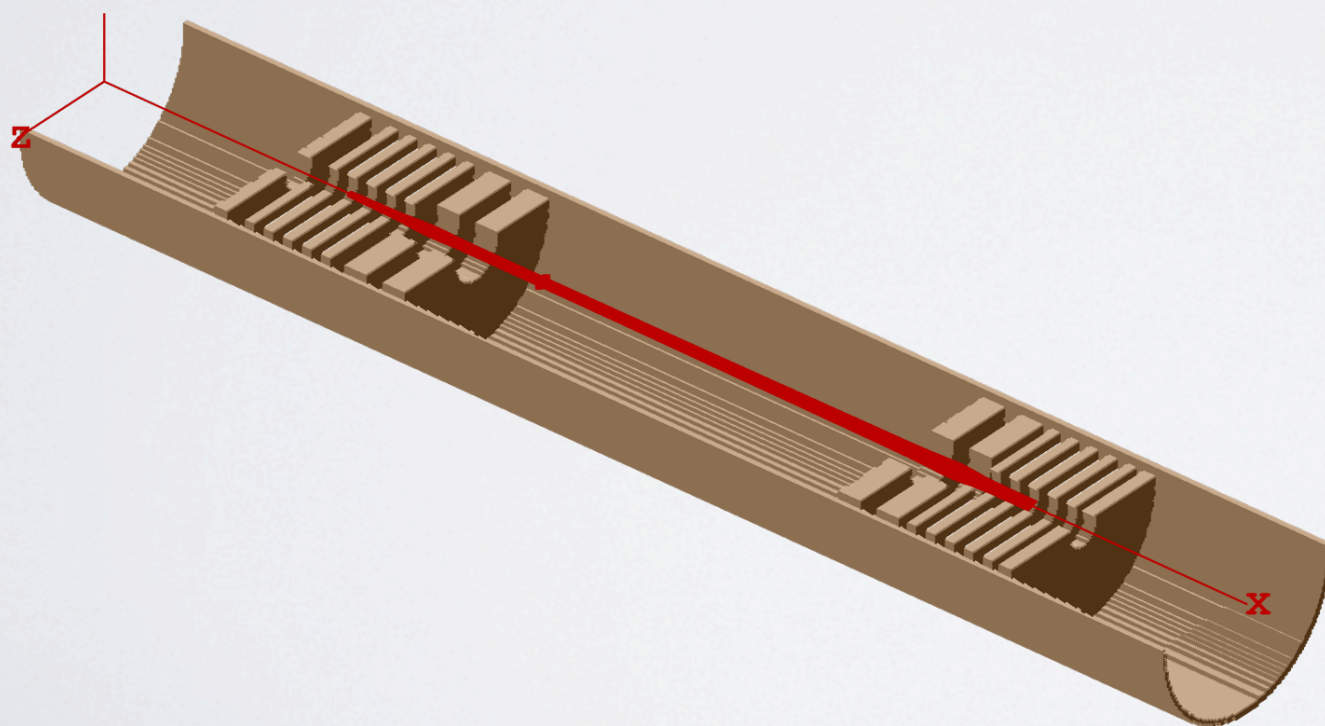








*Potential Contours*



*Simulated Ions in Trap*

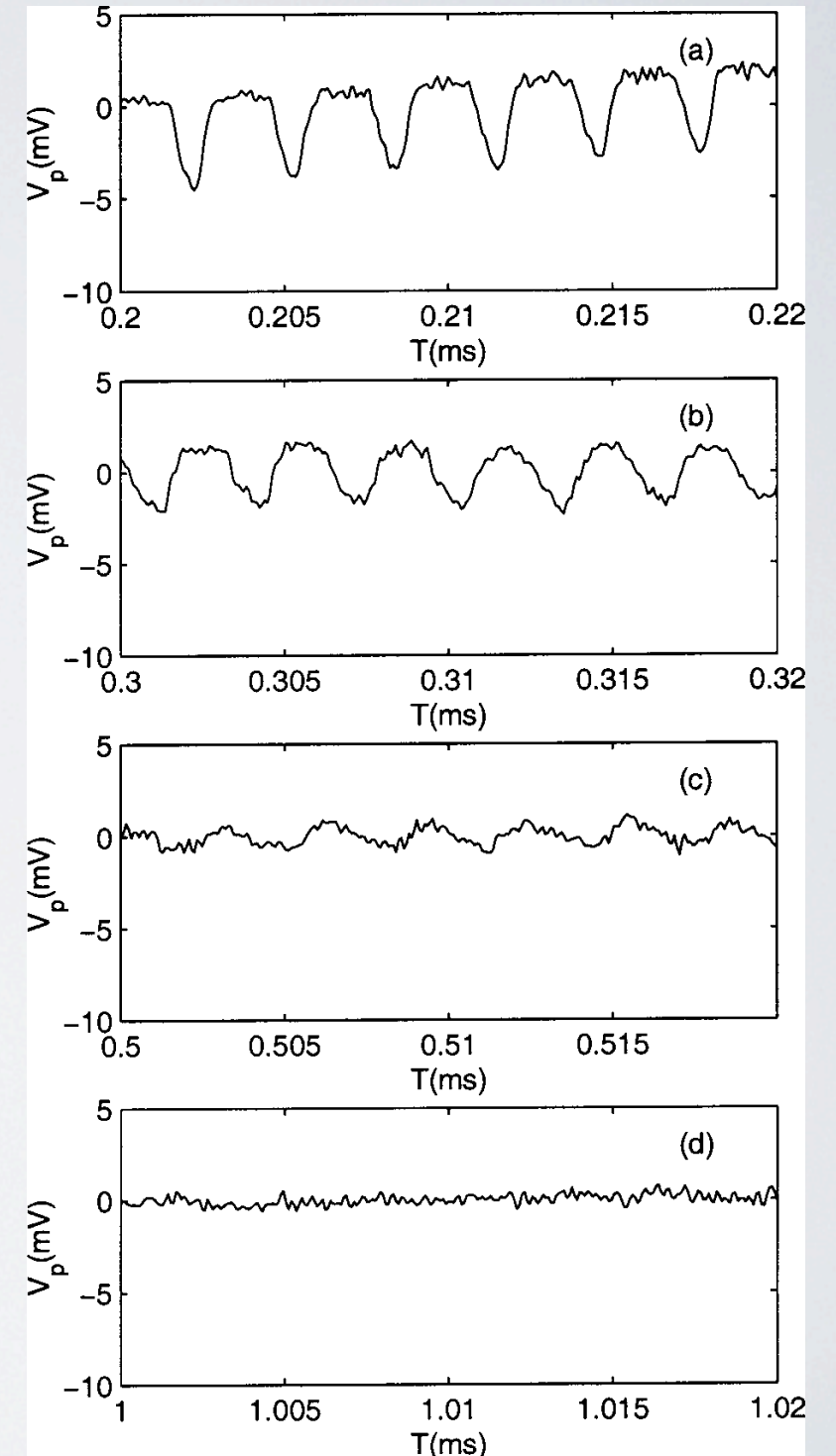


# Ion Behavior In the Trap

From simple arguments the width of the ion cloud in the trap should increase as a function of the oscillation number (not all ion have the exact same energy).

$$W_n = (W_0^2 + n^2 \Delta T^2)^{1/2}$$

*Signal in pickup electrode for different times after injection.*





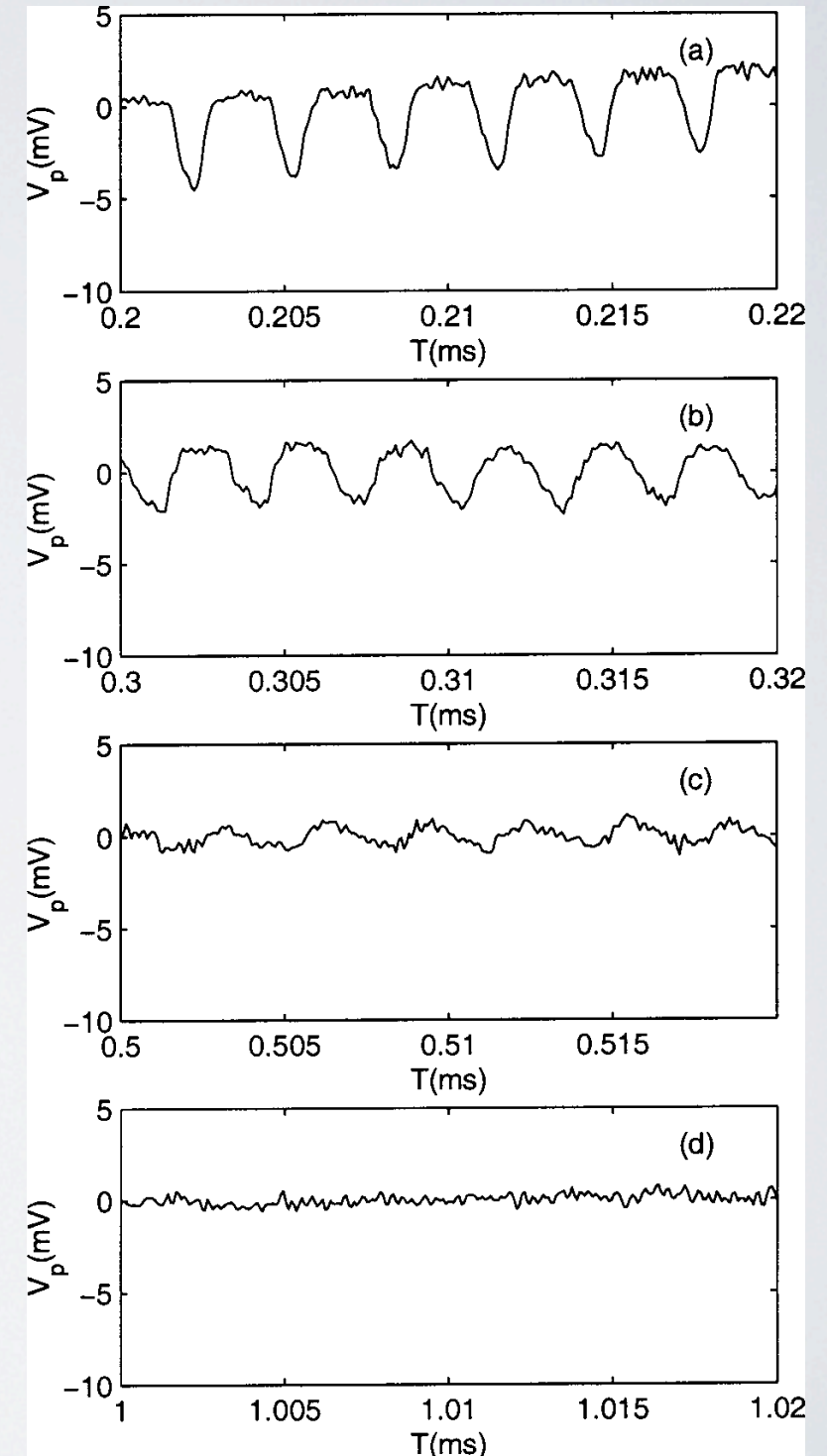
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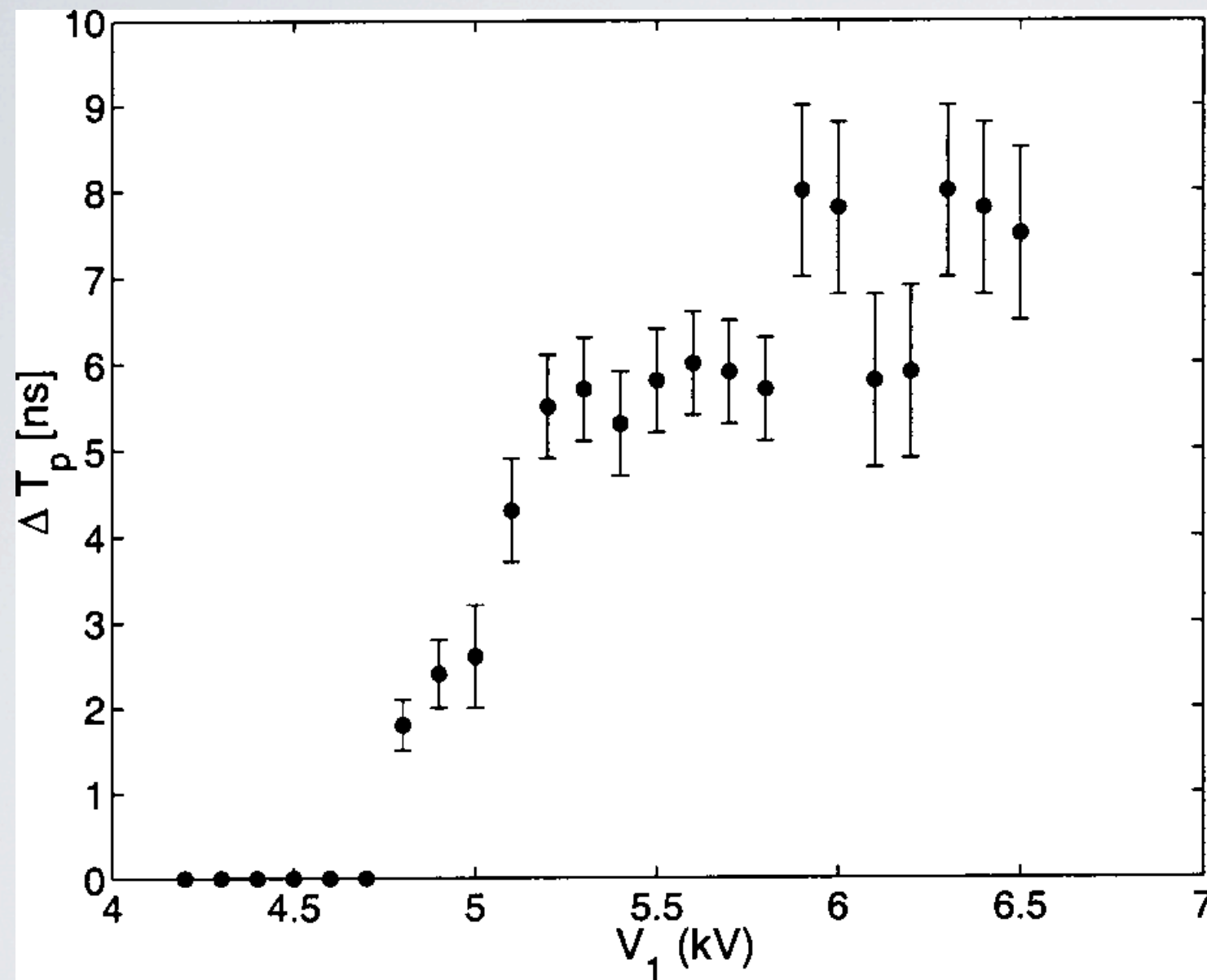
*Signal in pickup electrode for different times after injection.*

Using the pickup, it is possible to measure the detuning coefficient for different values of the (outer) electrode potential.



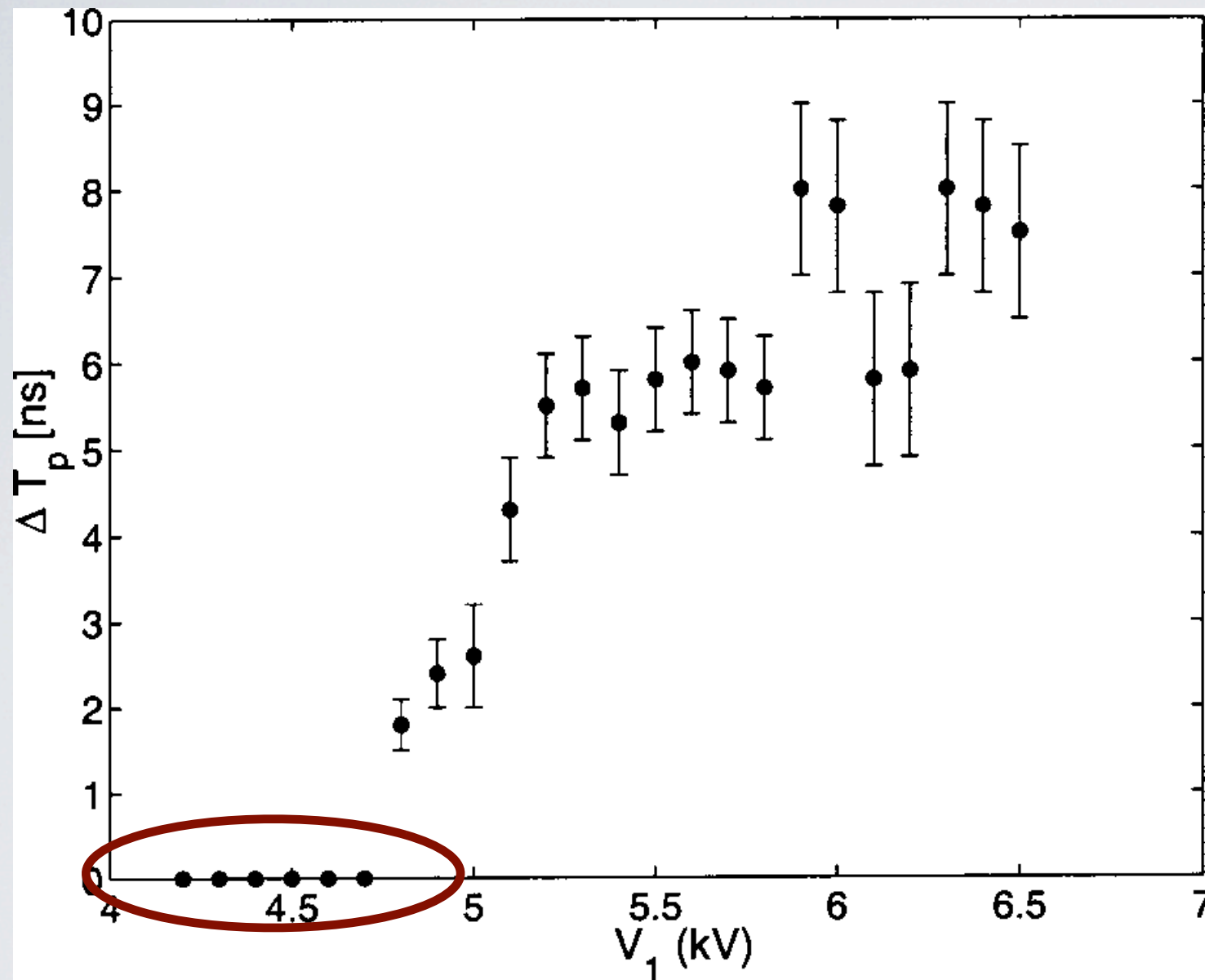


# Surprise...

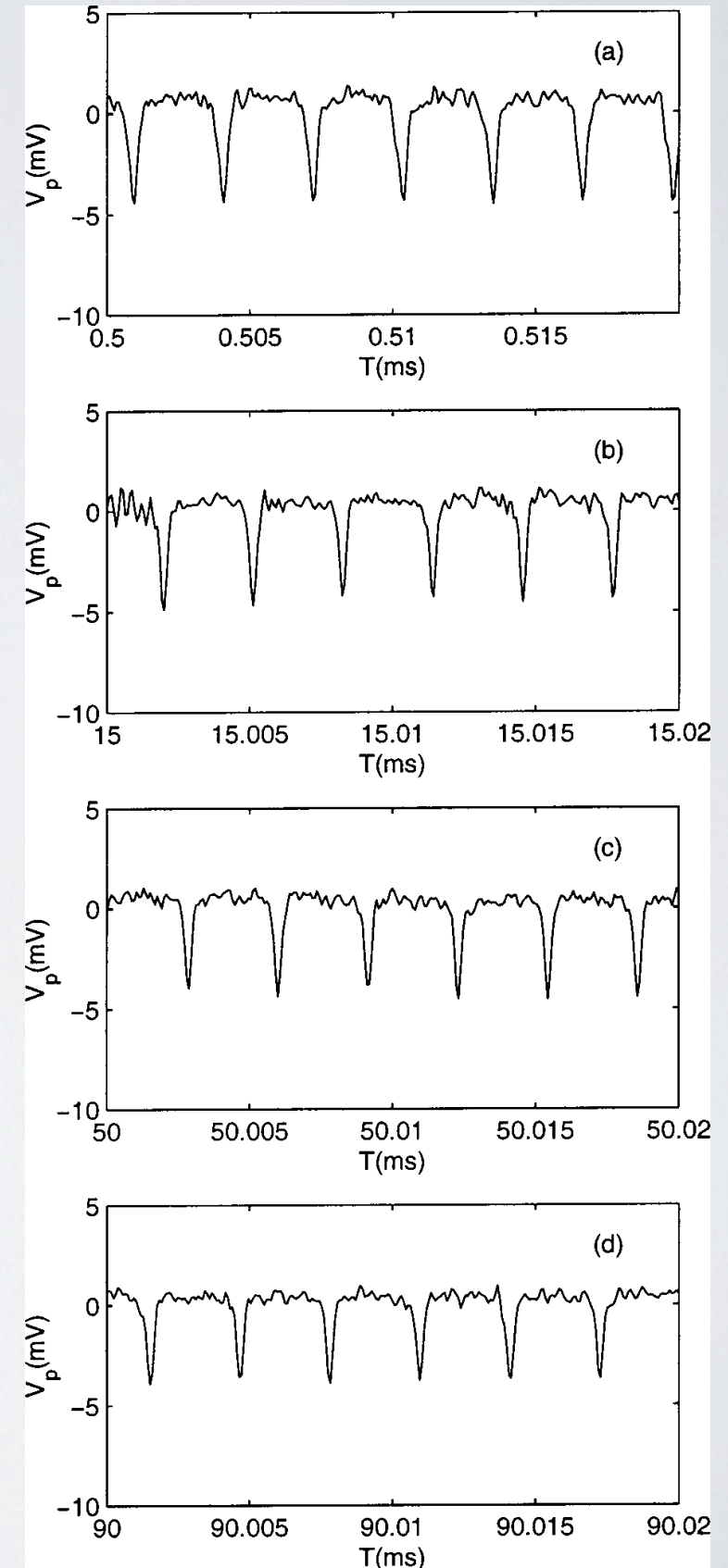




# Surprise...



For some values of the potential there is no dispersion!





# Why?

$$V(x) = \begin{cases} 0 & |x| \leq L/2 \\ F(|x| - L/2) & |x| > L/2 \end{cases} \quad \text{1D model for the potential in the trap}$$

$$T = 4 \left( \frac{L}{2\nu} + \frac{m\nu}{qF} \right) \quad \text{Oscillation period for ion with initial velocity } \nu$$

$$\frac{dT}{d\nu} = \frac{4m}{qF} - \frac{2L}{\nu^2} = 0 \quad \text{Extremum (minimum) condition}$$

For  $V_i < 4.75$  kV (empirically, for 4.2 keV

Ar<sup>+</sup> ions)  $dT/d\nu > 0$



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Handwaving explanation (but can be  
"proved" analytically)



# Why?

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Higher energy ions spend longer in the mirror region. On the way back they speed up the lower energy ions and get slowed by them.



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$$T = 4 \left( \frac{m}{2} \right)^{1/2} \int_{-L/2}^{L/2} \sqrt{F(x) - E} dx \quad \text{ion with}$$

$$\frac{dT}{d\nu} = \frac{4n}{qF} \quad \text{condition}$$

If you were an accelerator physicist  
you would call this:  
“Negative Mass Instability”

For  $V_1 < 4.75$  kV (empirically, for 4.2 keV  
Ar<sup>+</sup> ions)  $dT/d\nu > 0$

Higher energy ions spend longer in the mirror  
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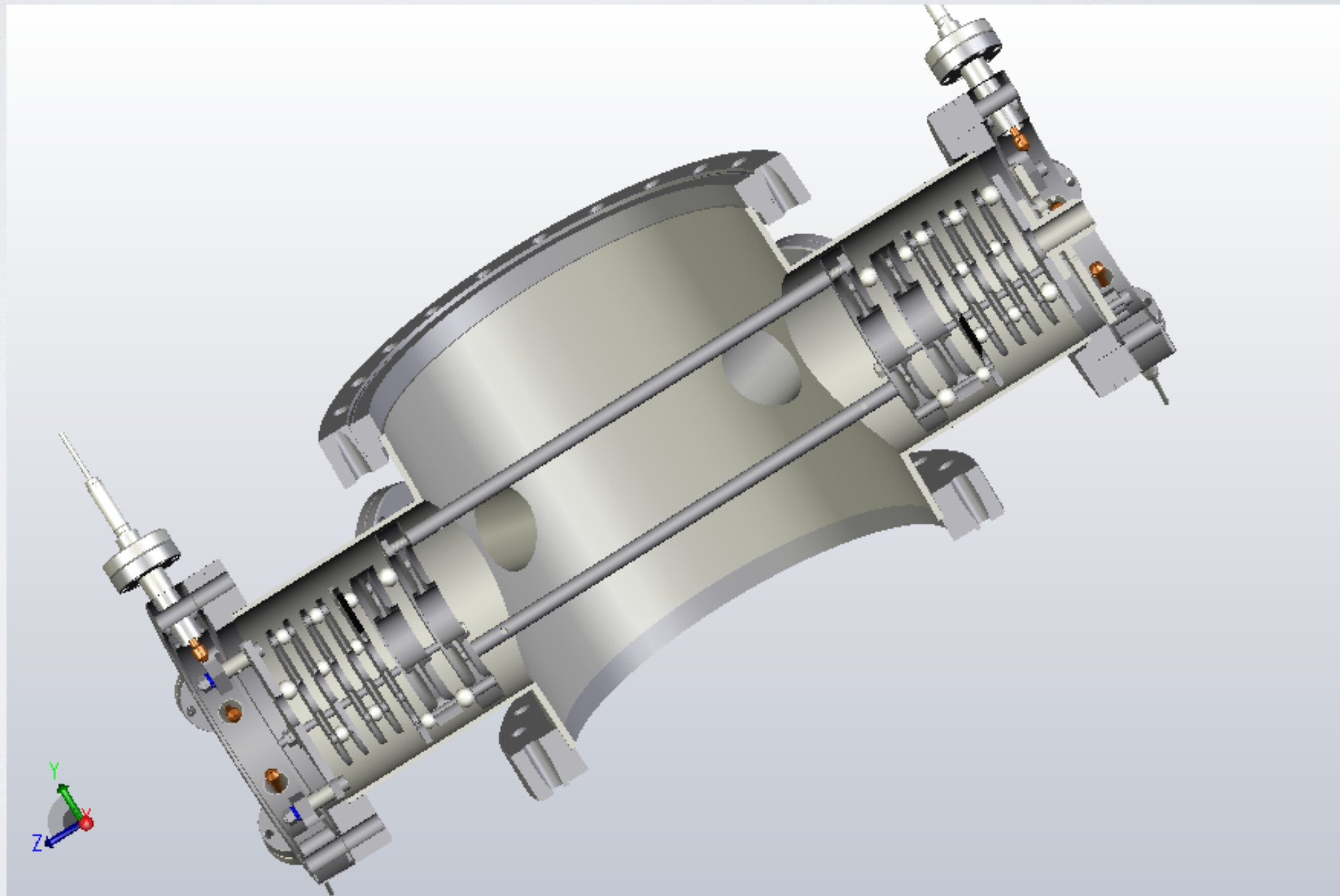
# POSSIBLE APPLICATIONS



# So what is it good for? (1)

*Mass*

*Spectroscopy*



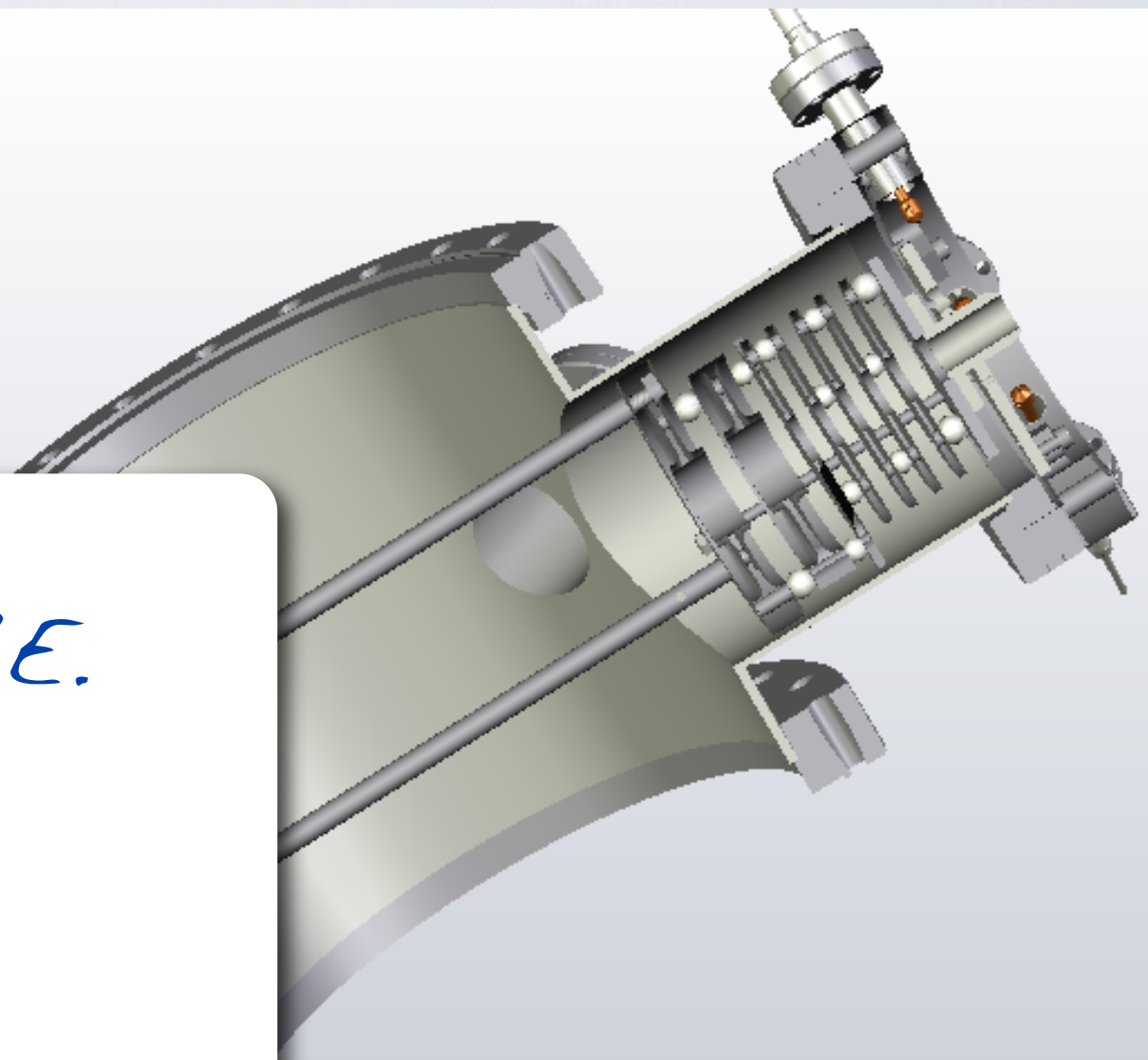


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*Trap selective in  $g/E$ .  
Different isotopes  
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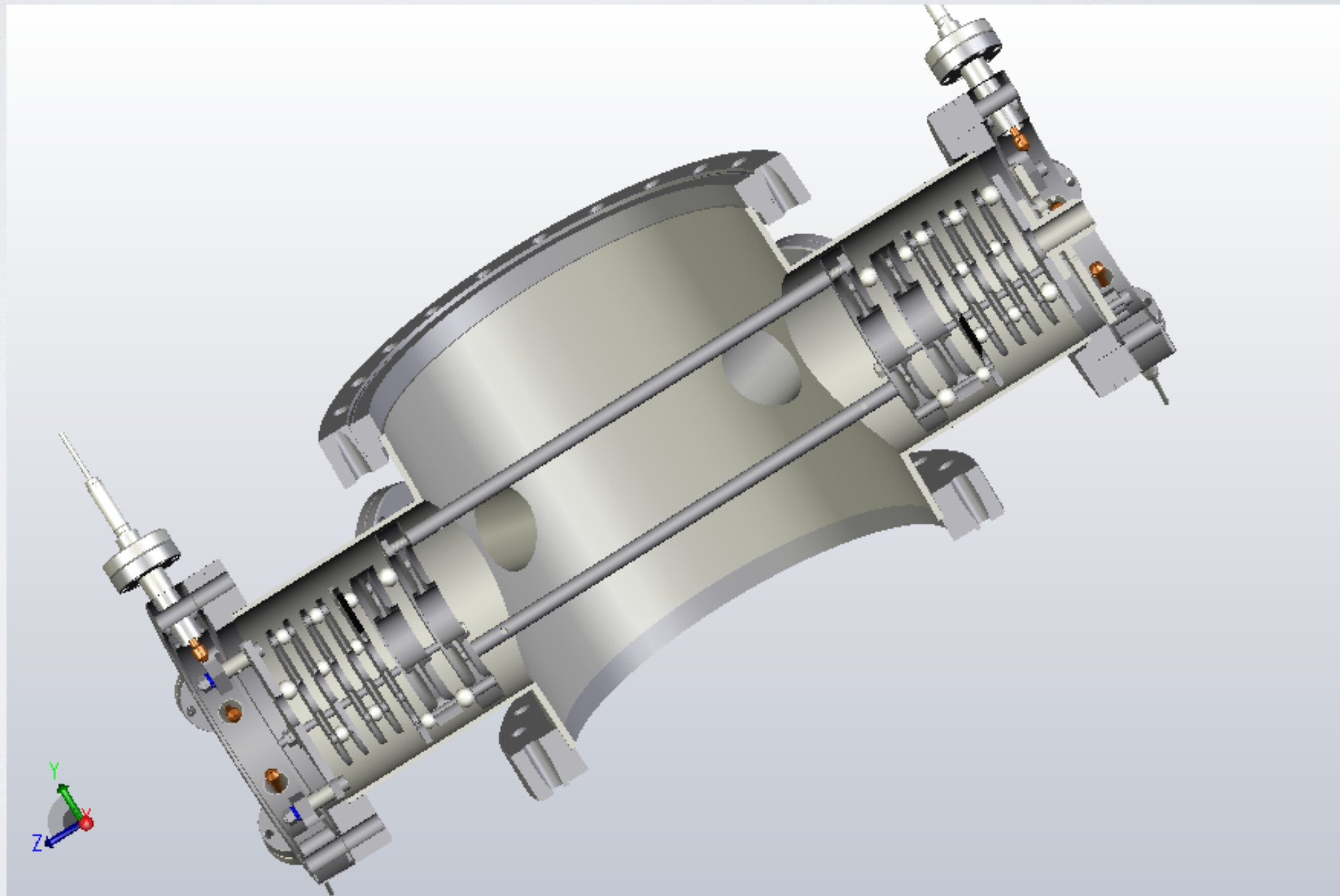




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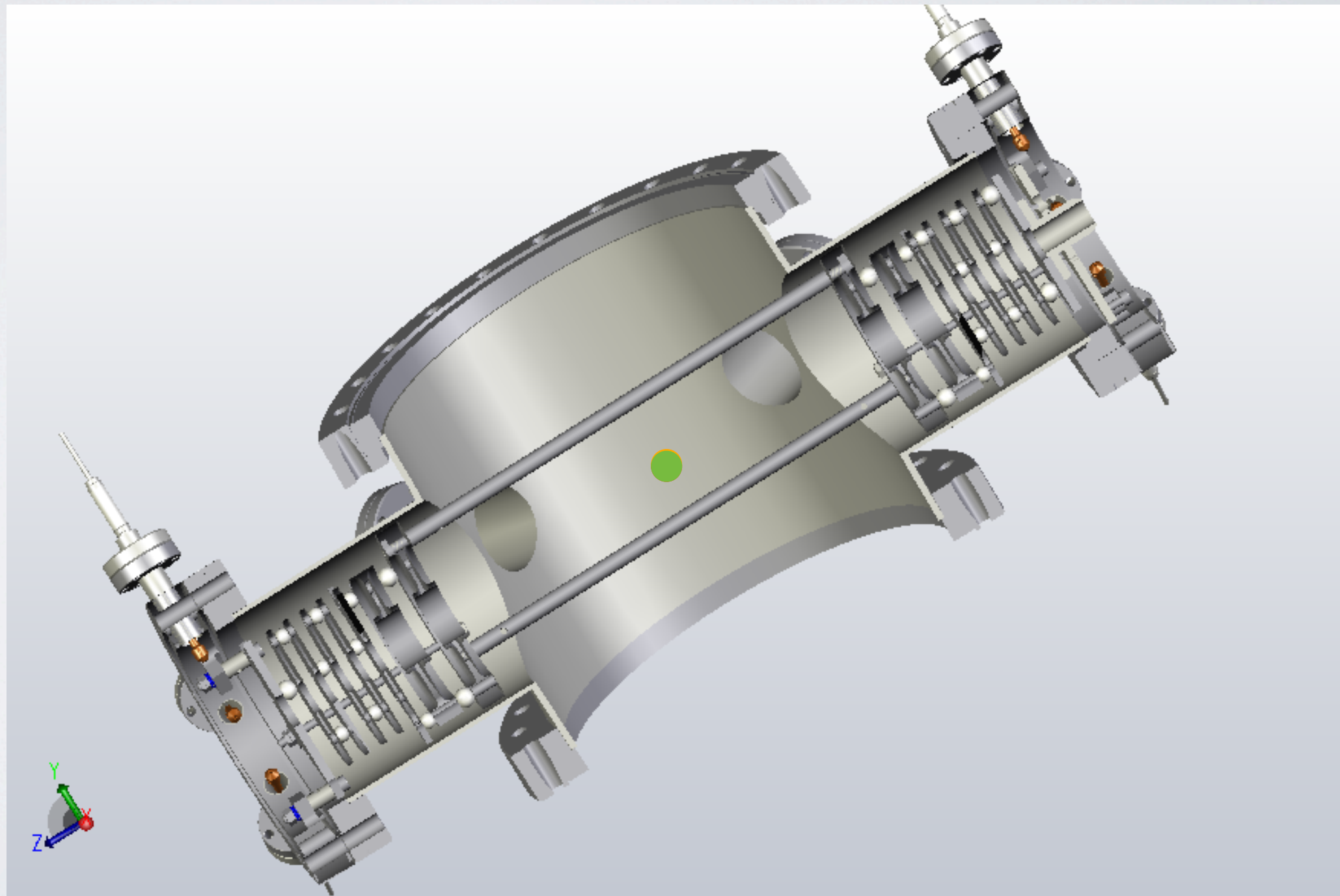




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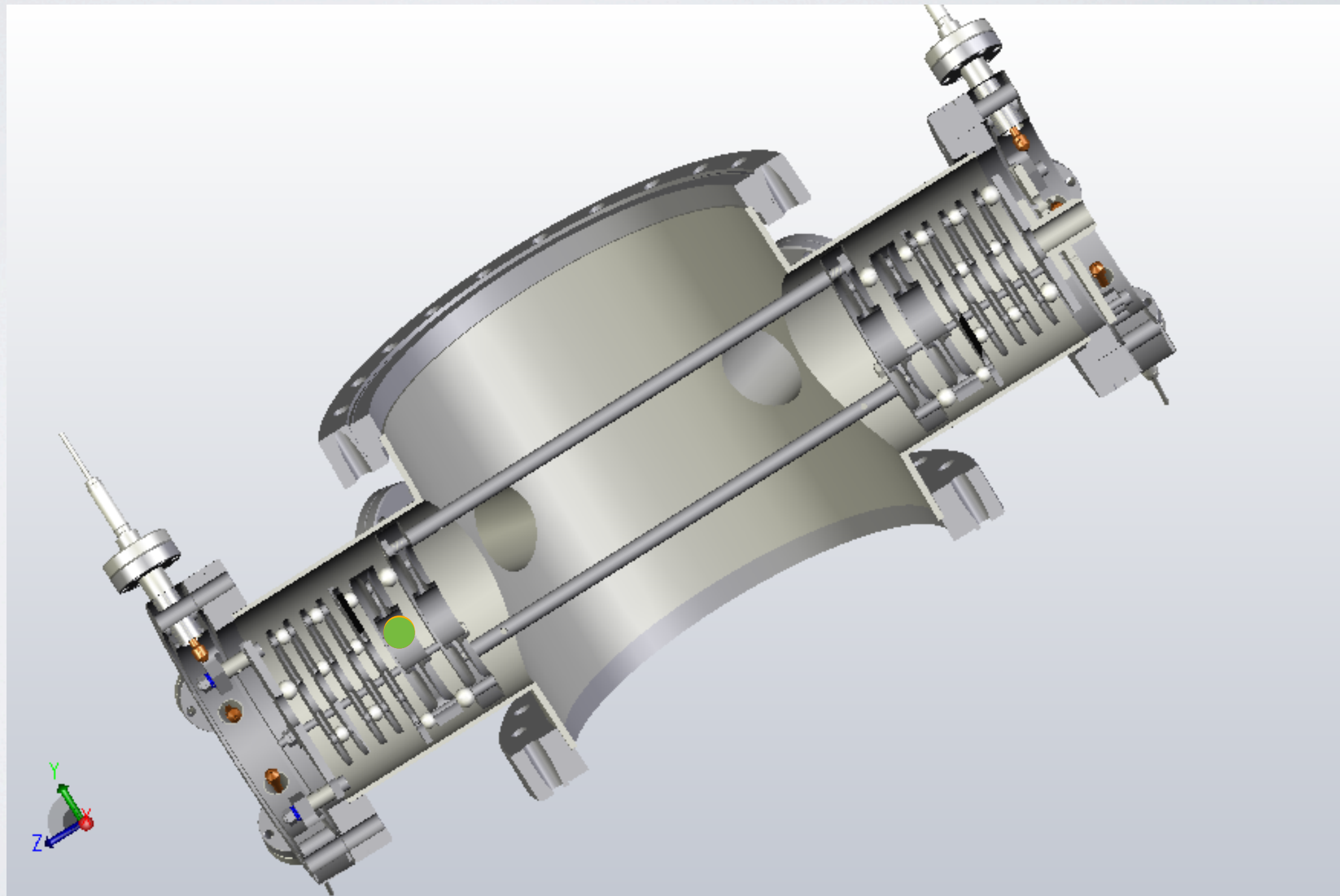




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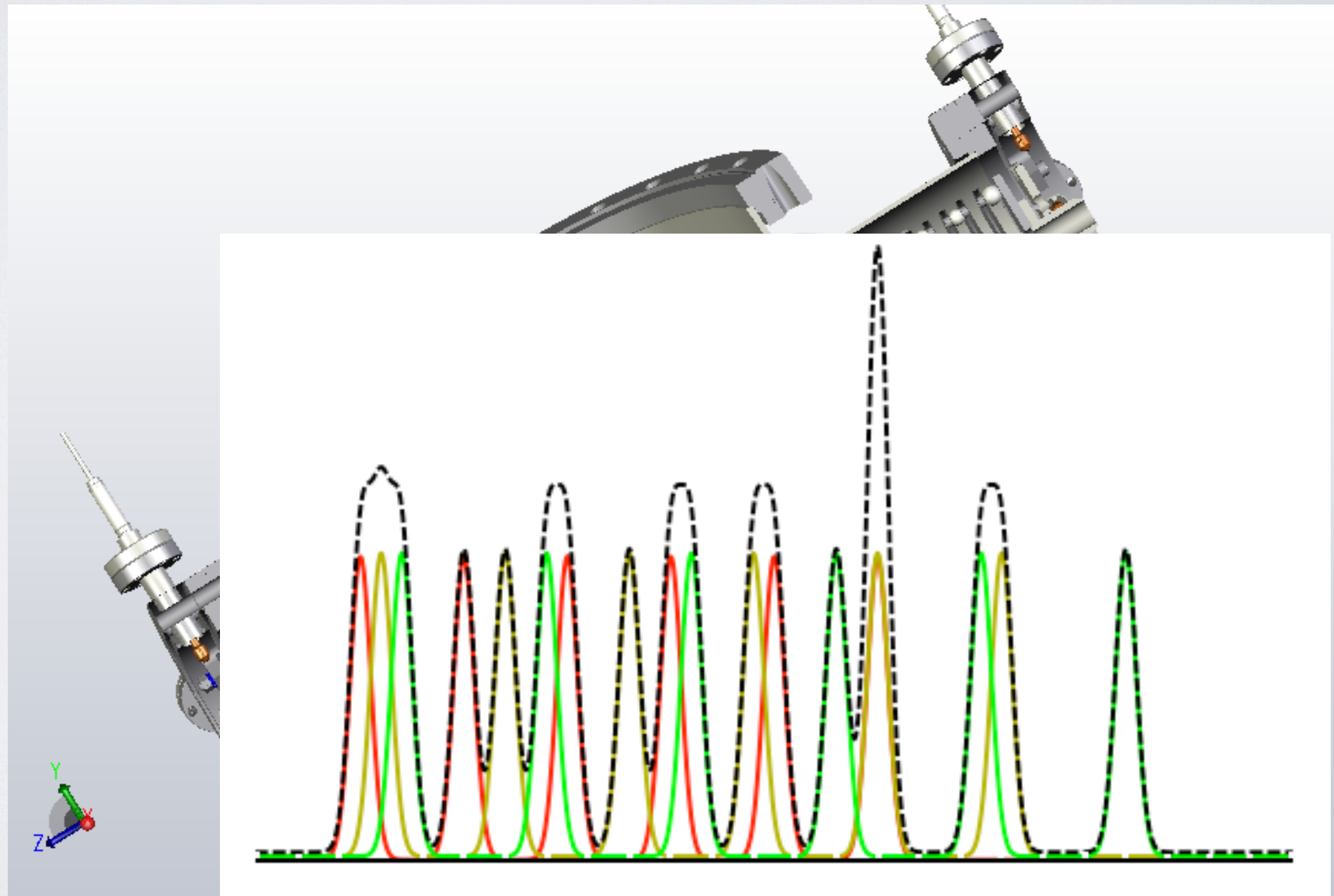




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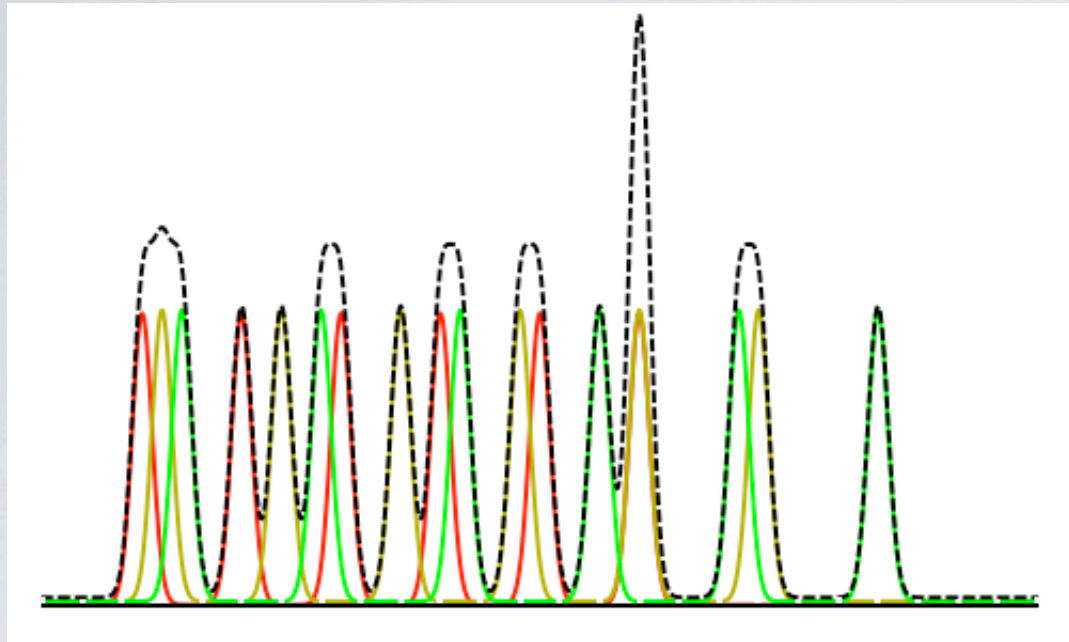
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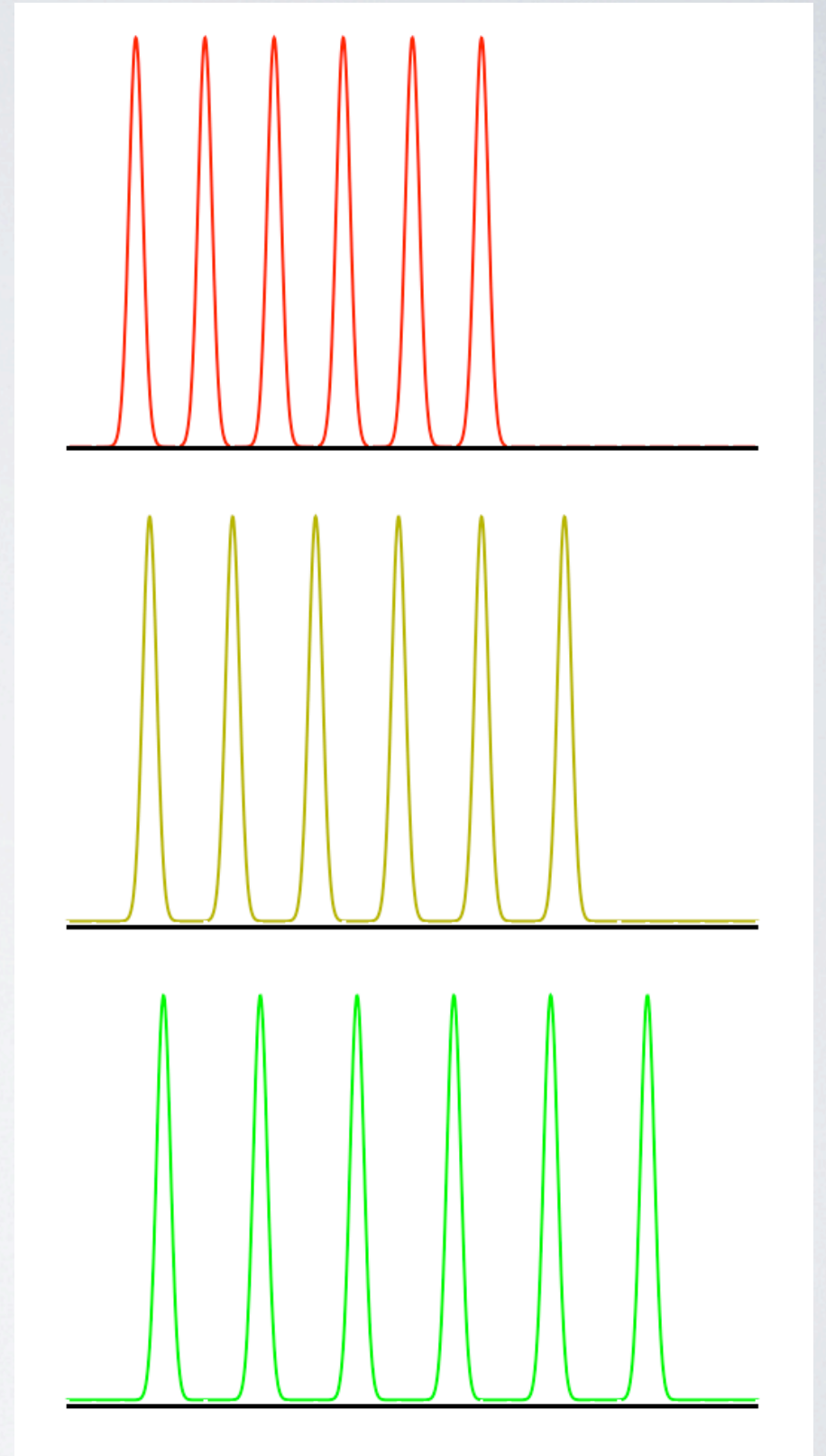




# Fourier Transform the pickup charge

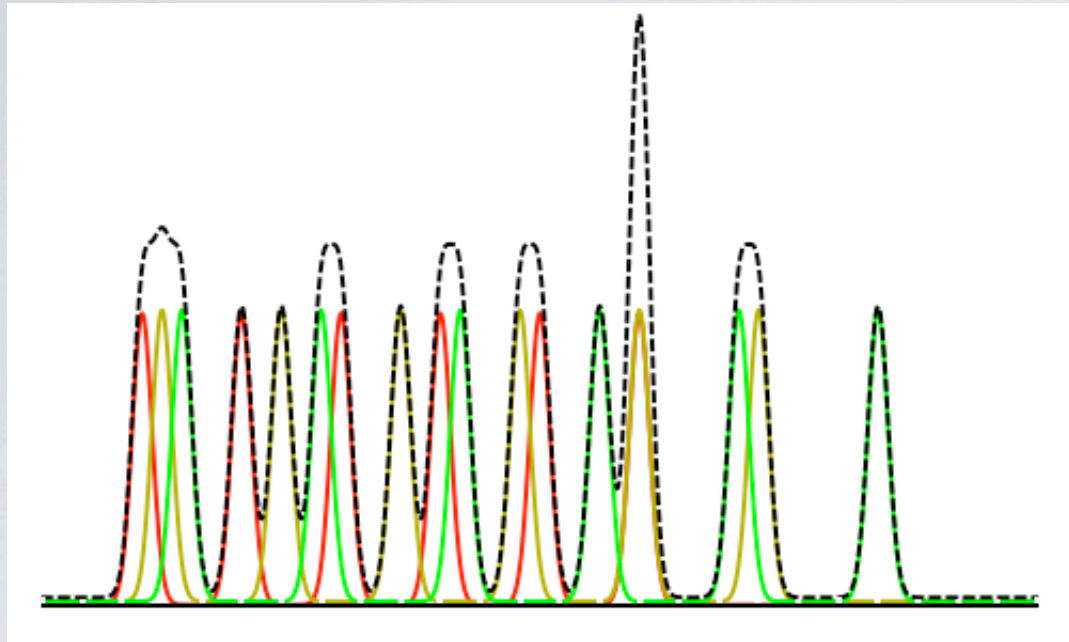


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Different isotopes  
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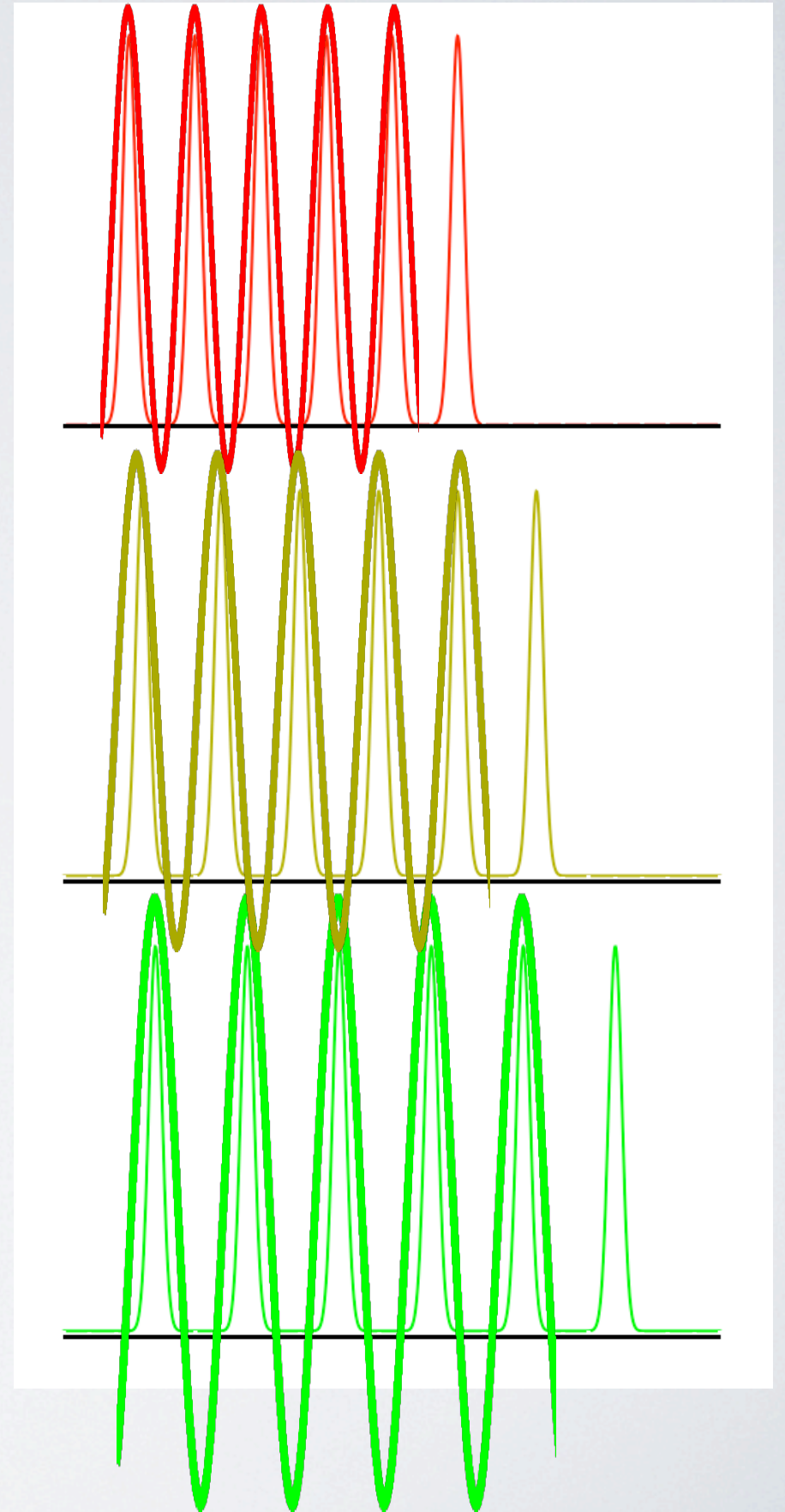




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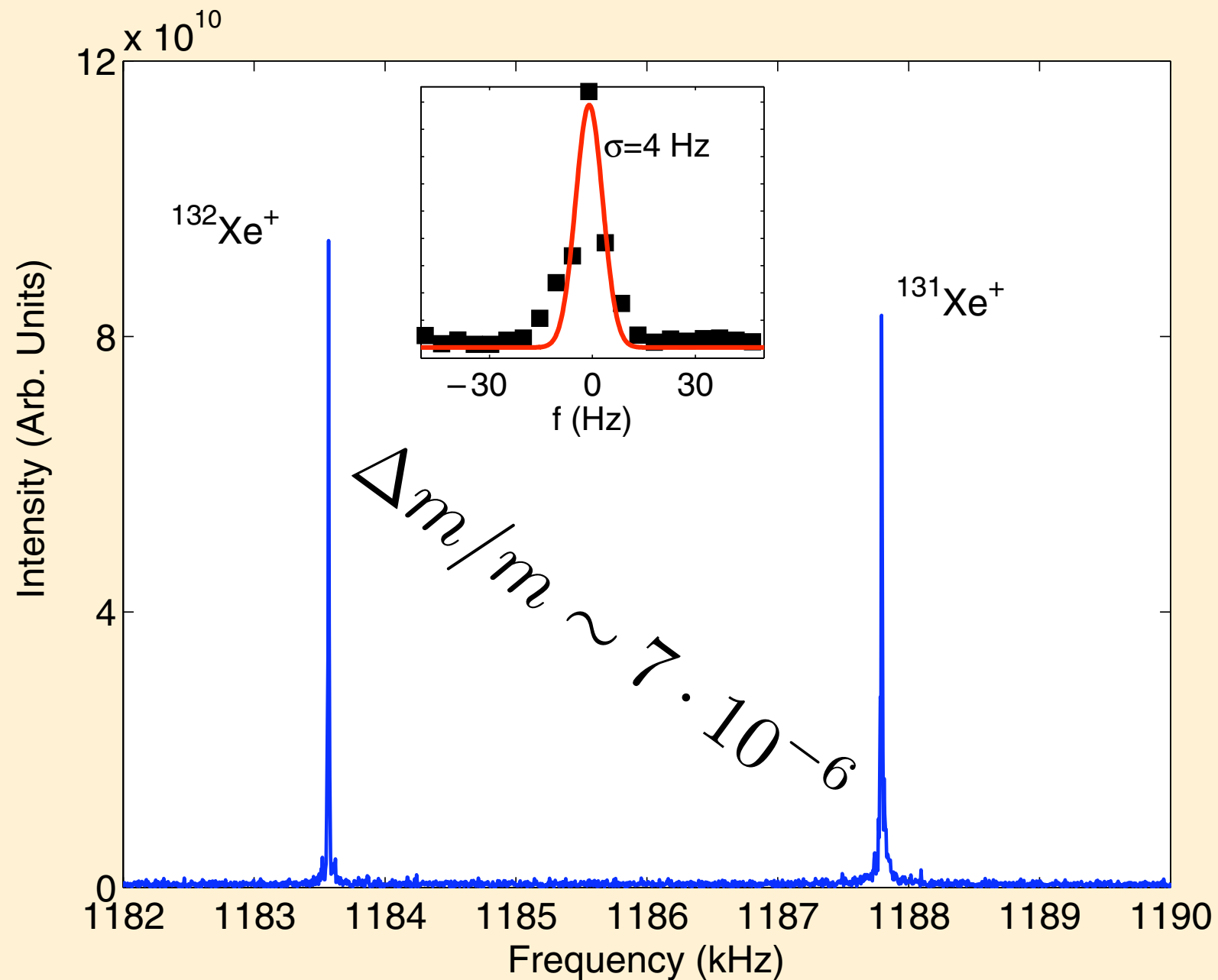


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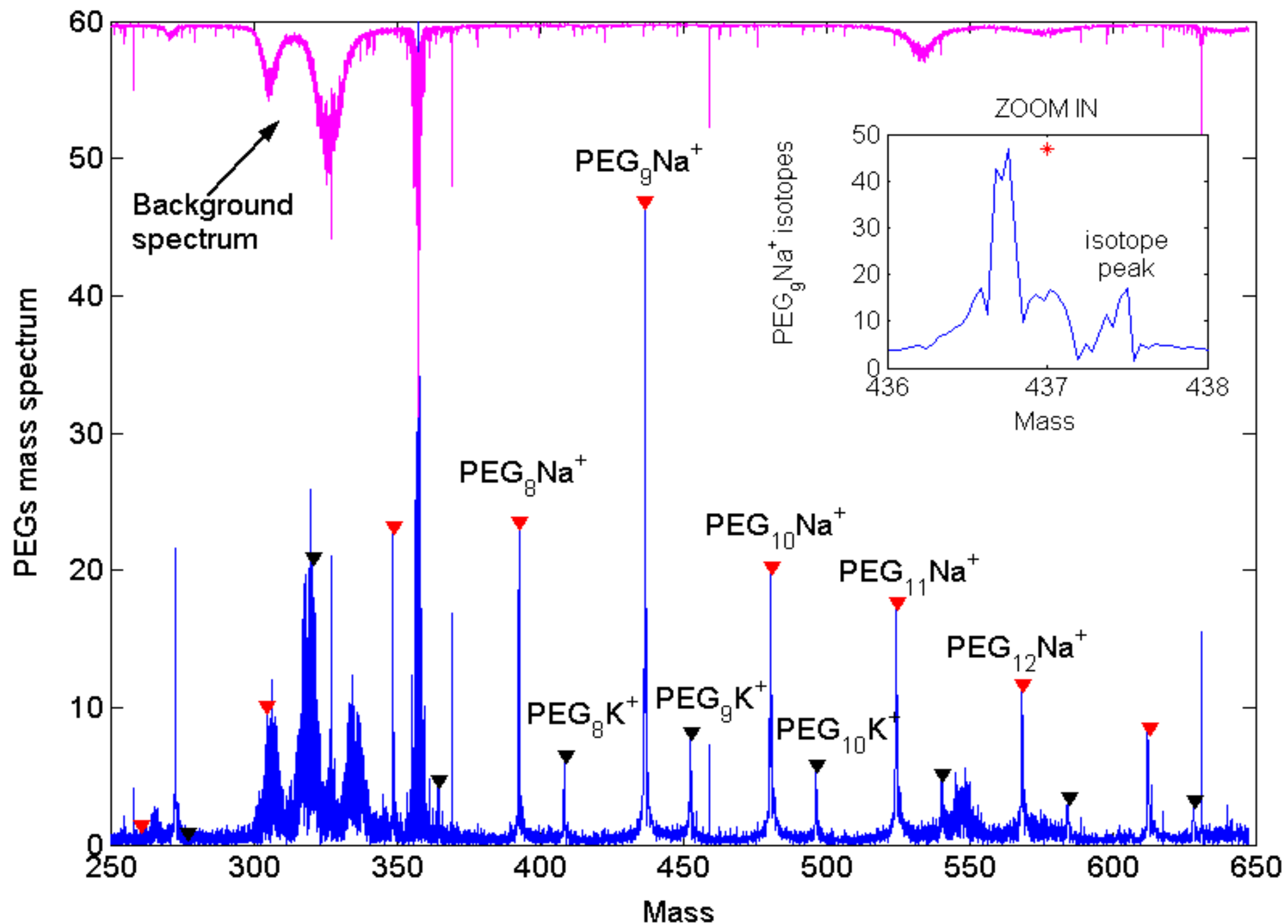


D. Strasser J. Phys. B: At. Mol. Opt. Phys. 36, 953 (2003)

Trap s  
Differ  
have a  
oscilla  
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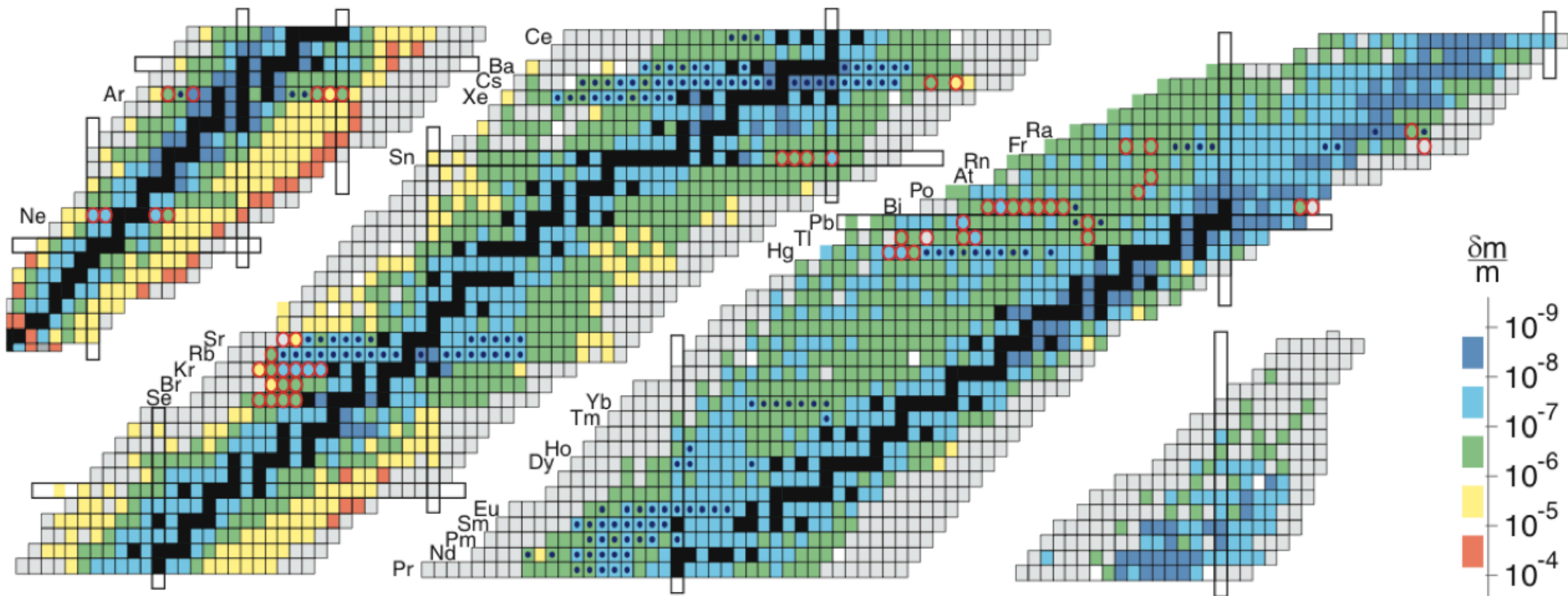


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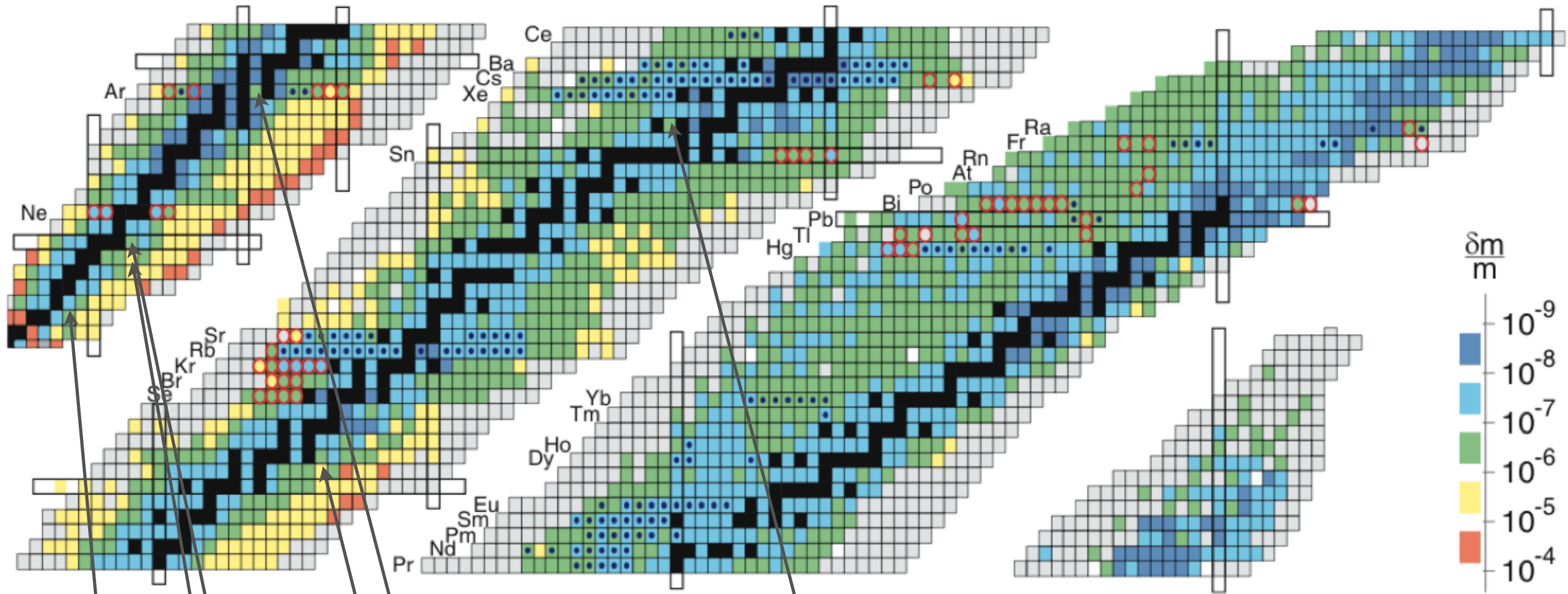


# World Mass Resolutions

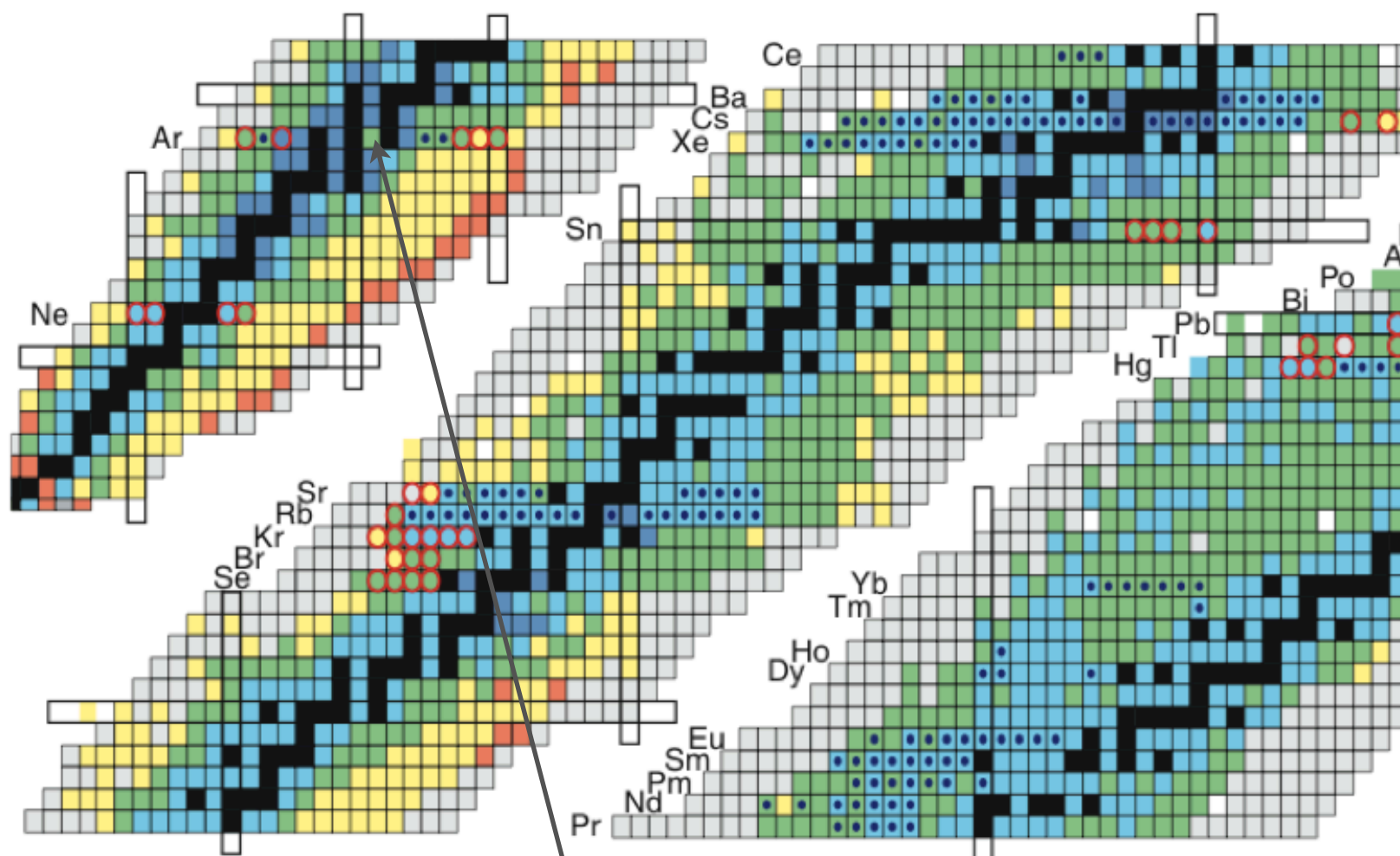


F. Herfurth et al., J. Phys. B: At. Mol. Opt. Phys. 36, 931 (2003)

# World Mass Resolutions





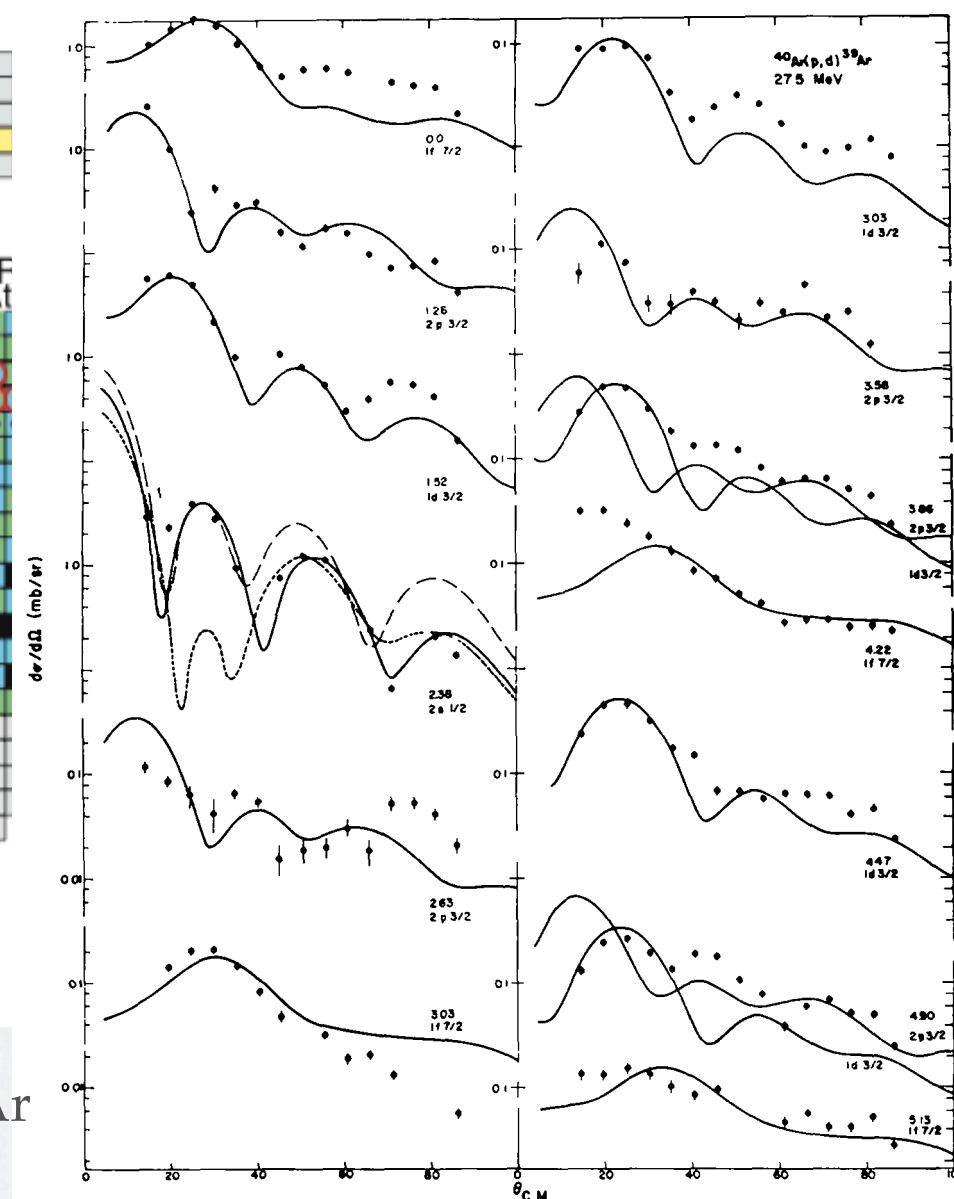


$^{39}\text{Ar}$

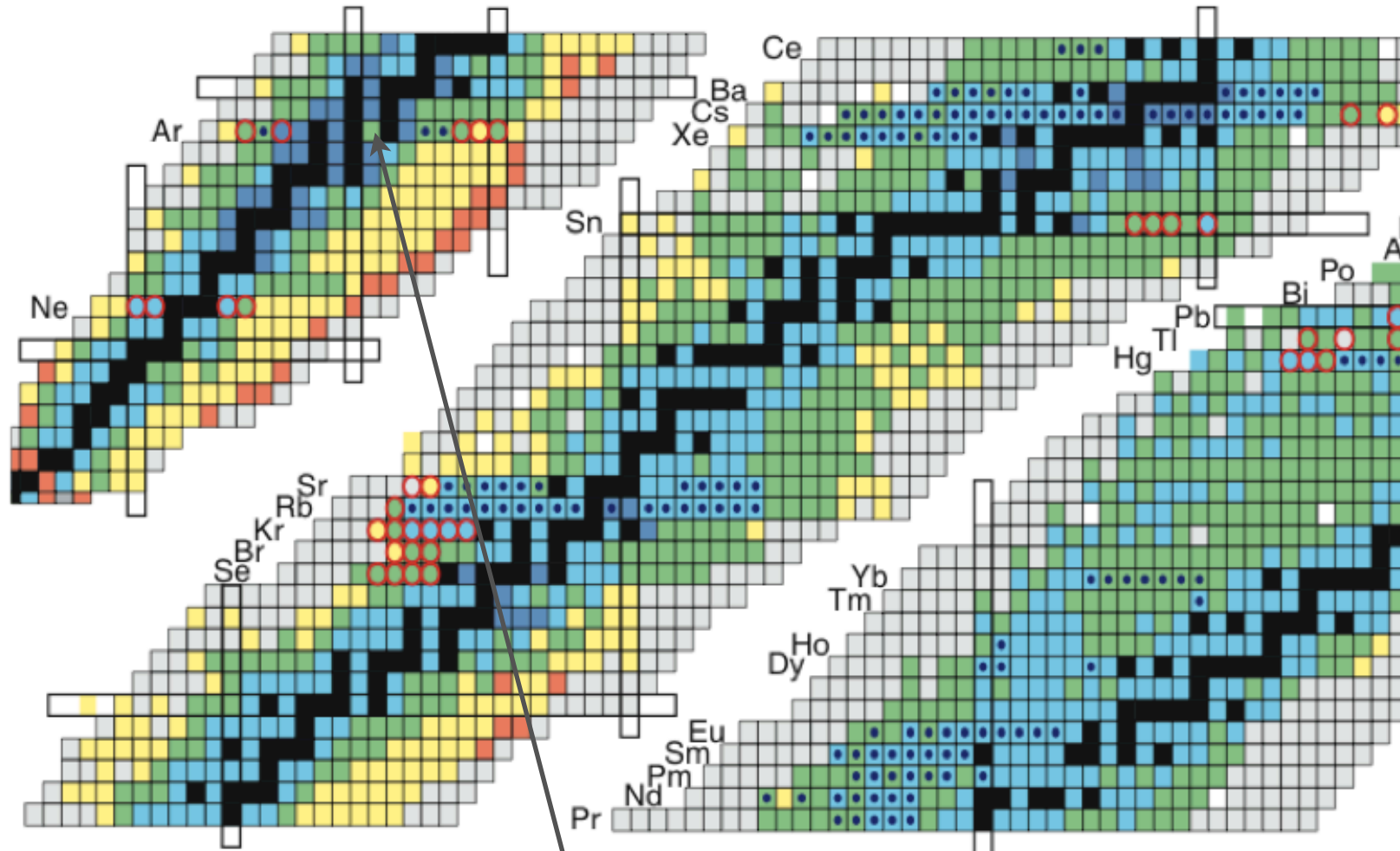
(269 yr)

- Produced via  $^{40}\text{Ar}(p,d)^{39}\text{Ar}$  with ~27 MeV protons.
- $^{40}\text{Ar}$  pretty much the easiest stuff to get.
- Long lived - easy transport.
- $^{38,40}\text{Ar}$  stable - perfect for calibration.
- All-in-all a good test case.

*F. Herfurth et al., J. Phys. B: At. Mol. Opt. Phys. 36, 931 (2003)*



*R. R. Johnson et al., Nucl Phys A108, 113 (1968)*

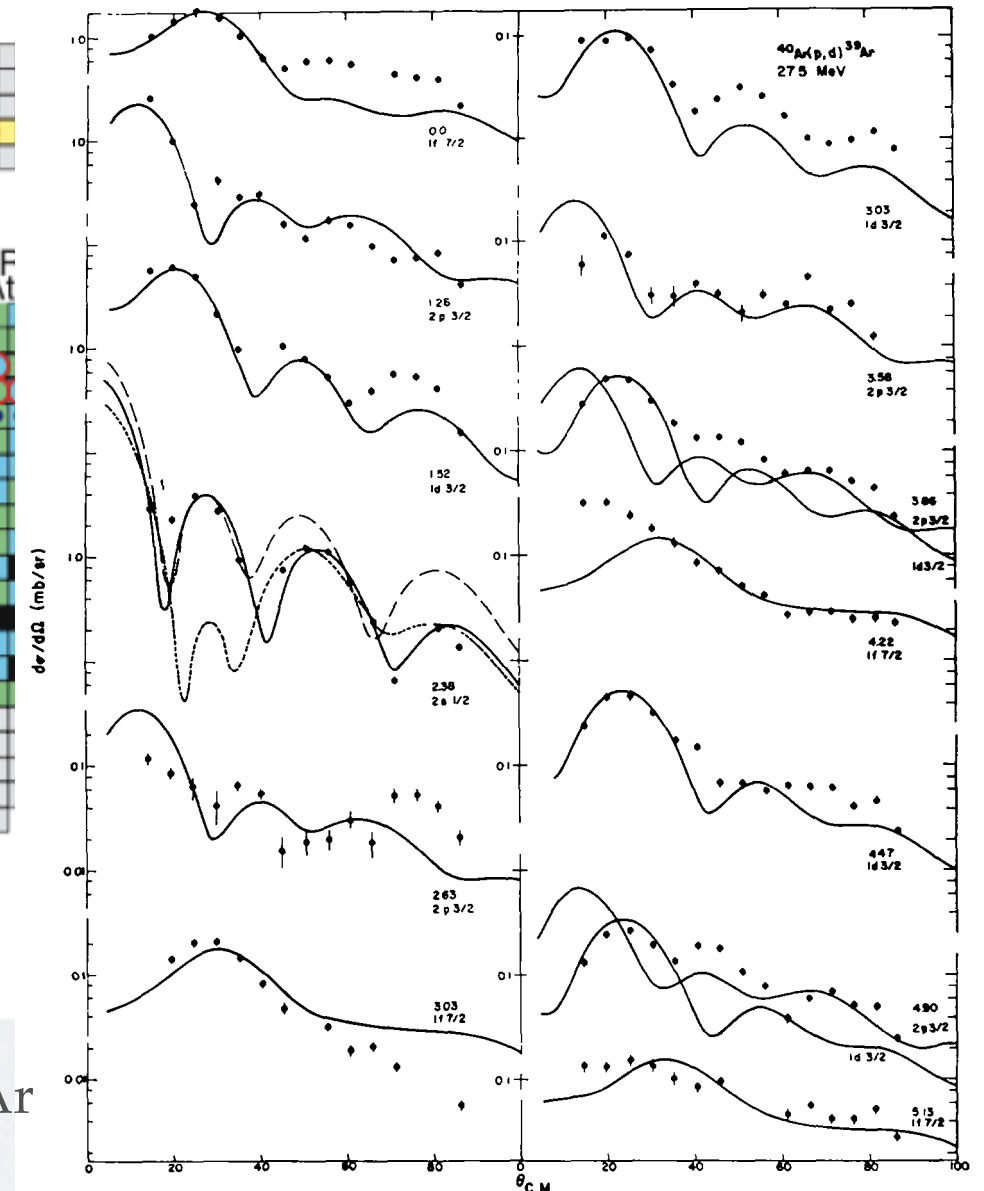


What about  
superheavy  
elements?

$^{39}\text{Ar}$   
(269 yr)

- Produced via  $^{40}\text{Ar}(p,d)^{39}\text{Ar}$  with ~27 MeV protons.
- $^{40}\text{Ar}$  pretty much the easiest stuff to get.
- Long lived - easy transport.
- $^{38,40}\text{Ar}$  stable - perfect for calibration.
- All-in-all a good test case.

*F. Herfurth et al., J. Phys. B: At. Mol. Opt. Phys. 36, 931 (2003)*

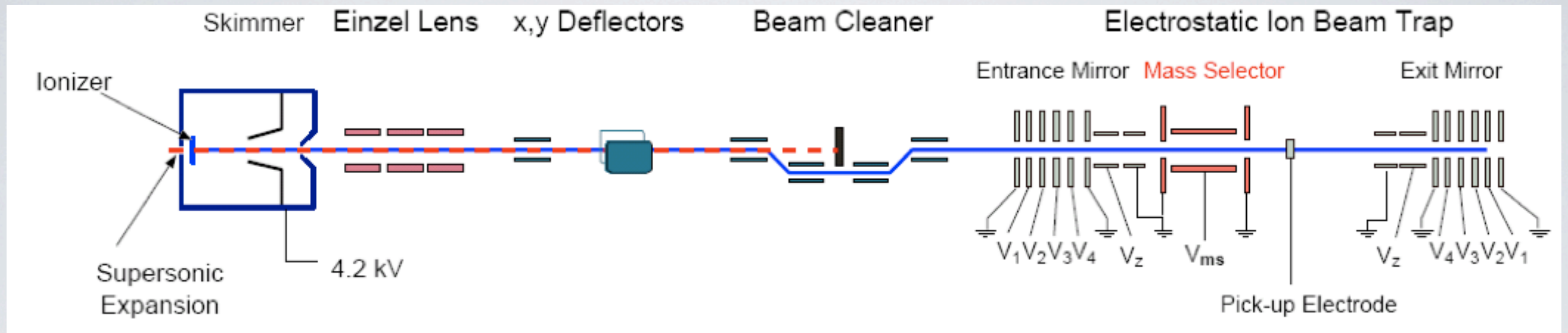


*R. R. Johnson et al., Nucl Phys A108, 113 (1968)*

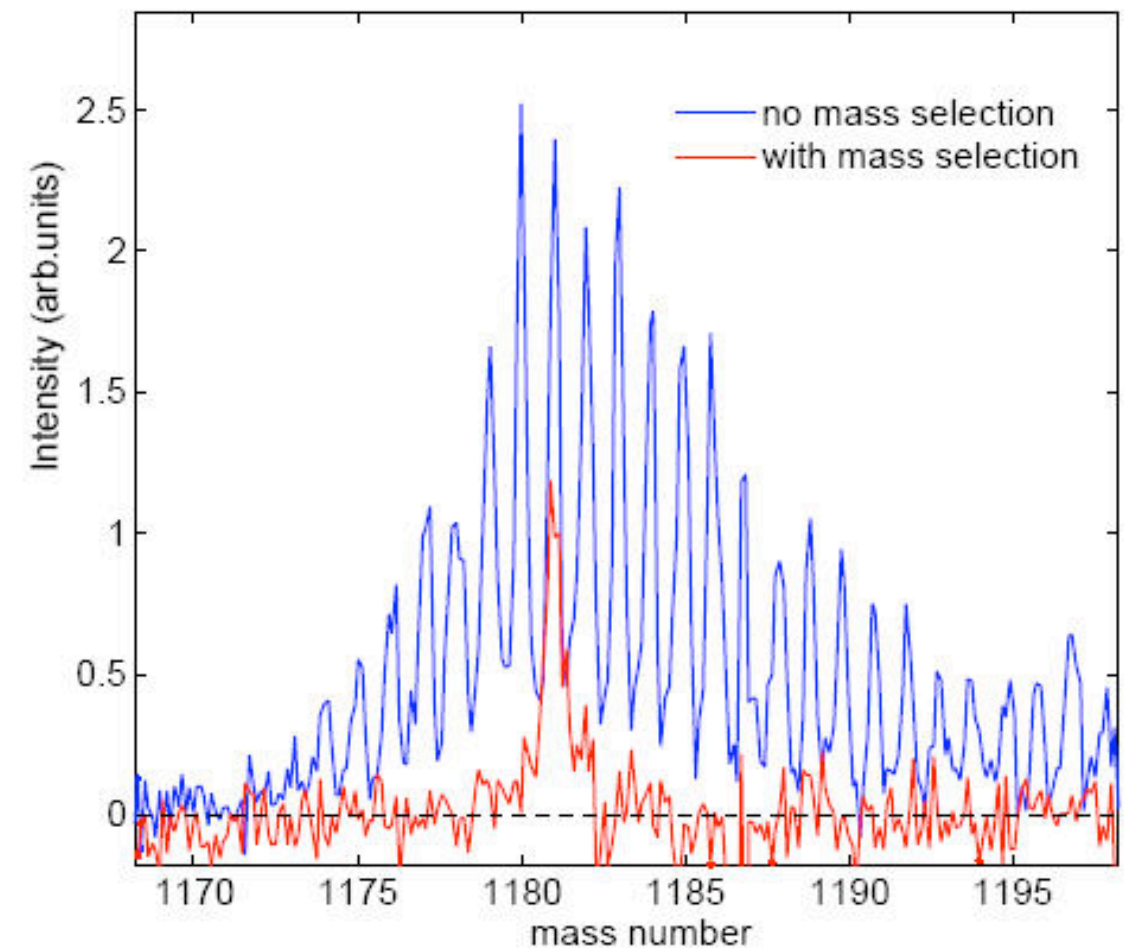


# So what is it good for? (2)

## Mass Selection



Apply RF pulse at correct frequency to “kick out” bunches with incorrect oscillation frequency.



# So what is it good for? (3)

## *$\beta$ decay studies*

$$\frac{d\Gamma}{dE_\beta d\Omega_\beta d\Omega_\nu} \propto \xi \left\{ 1 + \textcolor{red}{a} \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \textcolor{violet}{b} \frac{m}{E_e} + c \left[ \frac{1}{3} \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} - \frac{(\vec{p}_e \cdot \vec{j})(\vec{p}_\nu \cdot \vec{j})}{E_e E_\nu} \right] \right. \\ \left. \left[ \frac{J(J+1) - 3 \langle (\vec{J} \cdot \vec{j})^2 \rangle}{J(2J-1)} \right] + \frac{\langle \vec{J} \rangle}{J} \cdot \left[ \textcolor{blue}{A} \frac{\vec{p}_e}{E_e} + \textcolor{green}{B} \frac{\vec{p}_\nu}{E_\nu} + \textcolor{yellow}{D} \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right\}$$

*$\beta$  decay rate*

Parameter	Observable	Sensitivity	SM Prediction
$\textcolor{red}{a}$	$\beta$ - $\nu$ (recoil) correlation	Tensor & Scalar terms	$\textcolor{red}{1}$ for pure Fermi $\textcolor{red}{-1/3}$ for pure GT or combination
$\textcolor{violet}{b}$ (Fierz term)	Comparison of $\beta^+$ to EC rate	SV/T/A interference	$0$
$\textcolor{blue}{A}$	$\beta$ asymmetry for polarized nuclei	Tensor, ST/VA Parity	Nucleus dependent
$\textcolor{green}{B}$	$\nu$ asymmetry (recoil) for polarized nuclei	Tensor, TA/ST/VA/SA/VT Parity	Nucleus dependent
$\textcolor{brown}{D}$	Triple product	ST/VA Interference TRI	$0$



Ion Trap for  $\beta$  decay - been done before (or in planning)

- WITCH - ISOLDE / CERN
- TITAN - TRIUMF
- LPC-Cean - GANIL

Ion Trap for  $\beta$  decay - been done before (or in planning)

- WITCH - ISOLDE / CERN
- TITAN - TRIUMF
- LPC-Cean - GANIL
- Complicated
- Superconducting
- Big / Stationary



Ion Trap for  $\beta$  decay - been done before (or in planning)

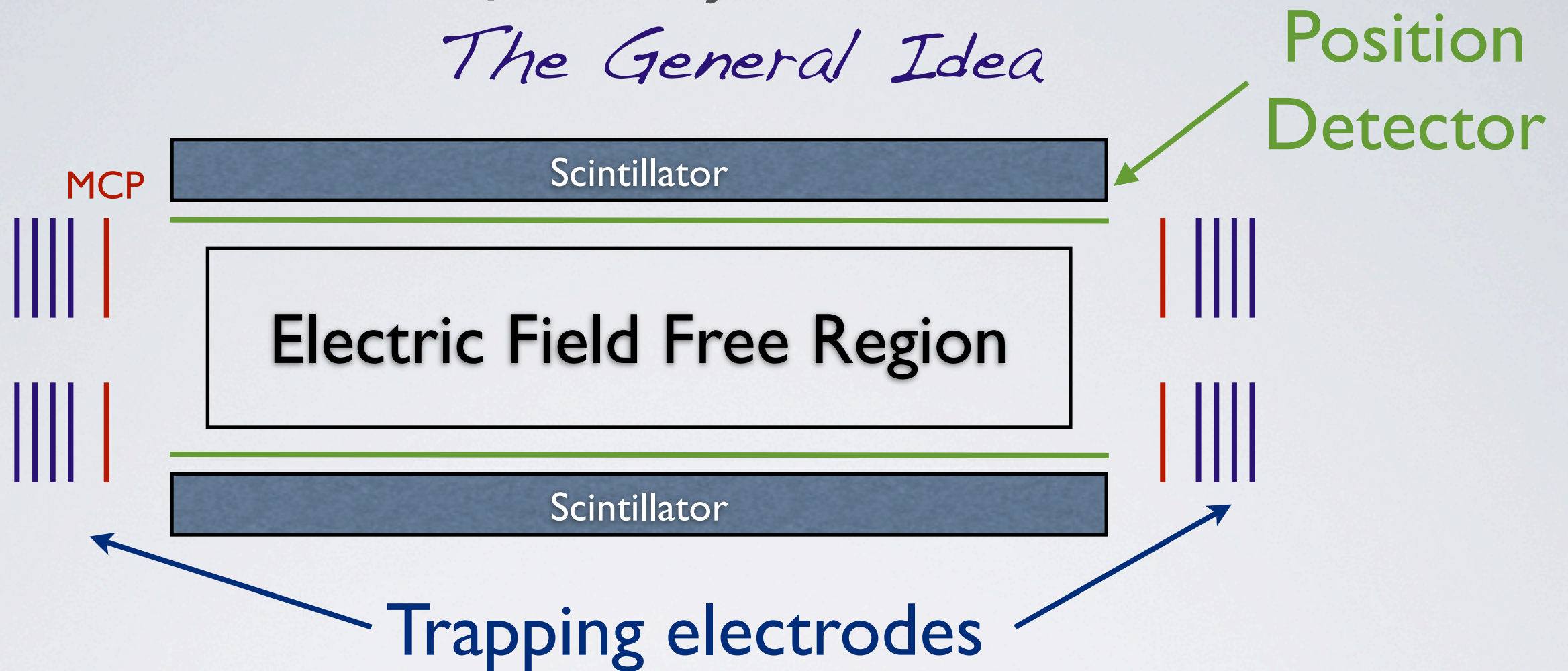
- WITCH - ISOLDE / CERN
- TITAN - TRIUMF
- LPC-Cean - GANIL
- Complicated
- Superconducting
- Big / Stationary

*Can we use an EIBT for this?*



# $\beta$ -Decay Studies

## *The General Idea*



- Recoil ion detected in MCP.
- $\beta$  detected in position detectors.
- Need bunch position for full reconstruction (multiple scattering of  $\beta$  in detectors).
- Large solid angle + kinematic focussing  $\rightarrow$  detection efficiency  $> 50\%$ .
- No need for electrostatic acceleration (ions at  $\sim \text{keV}$ ). Decay in field free region.

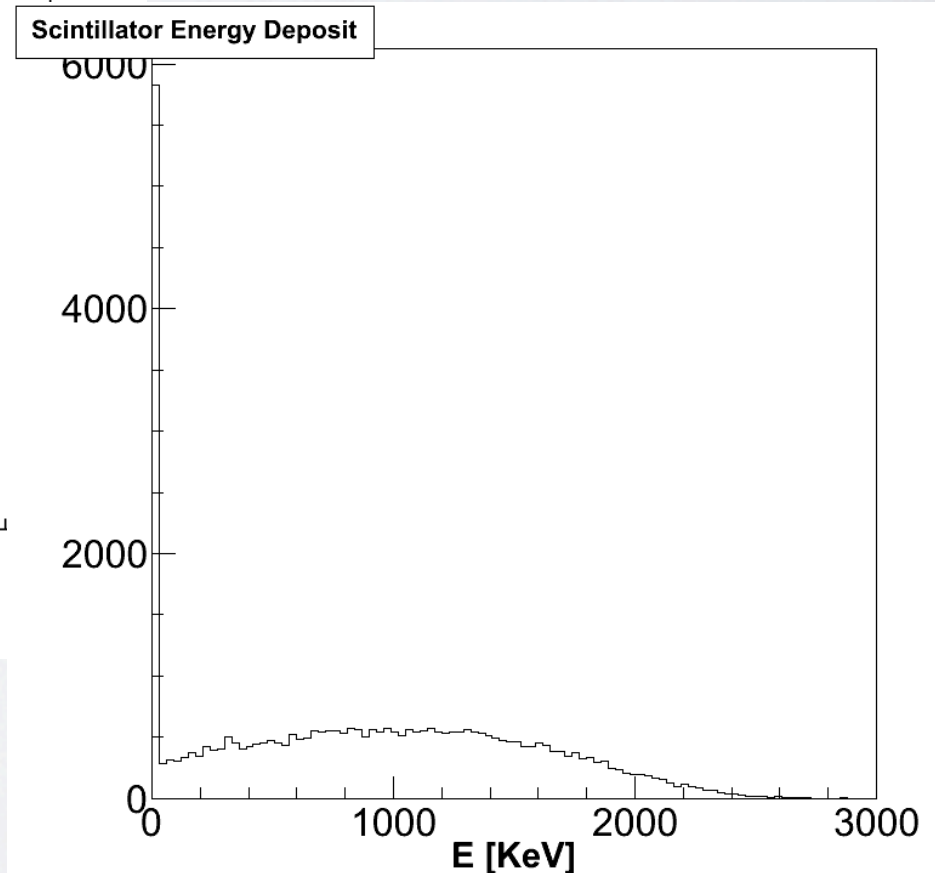
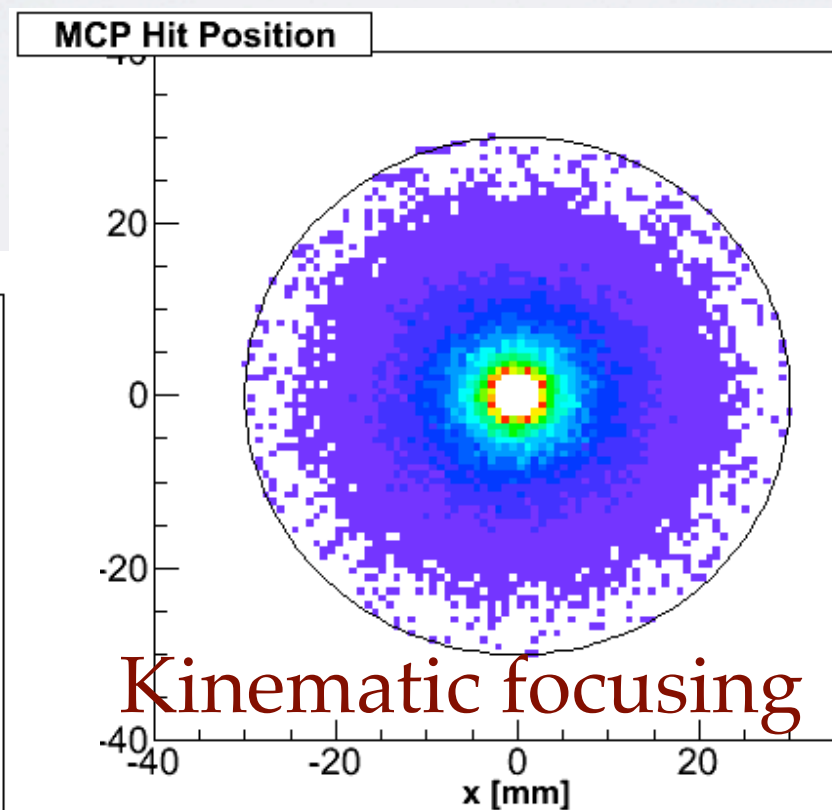
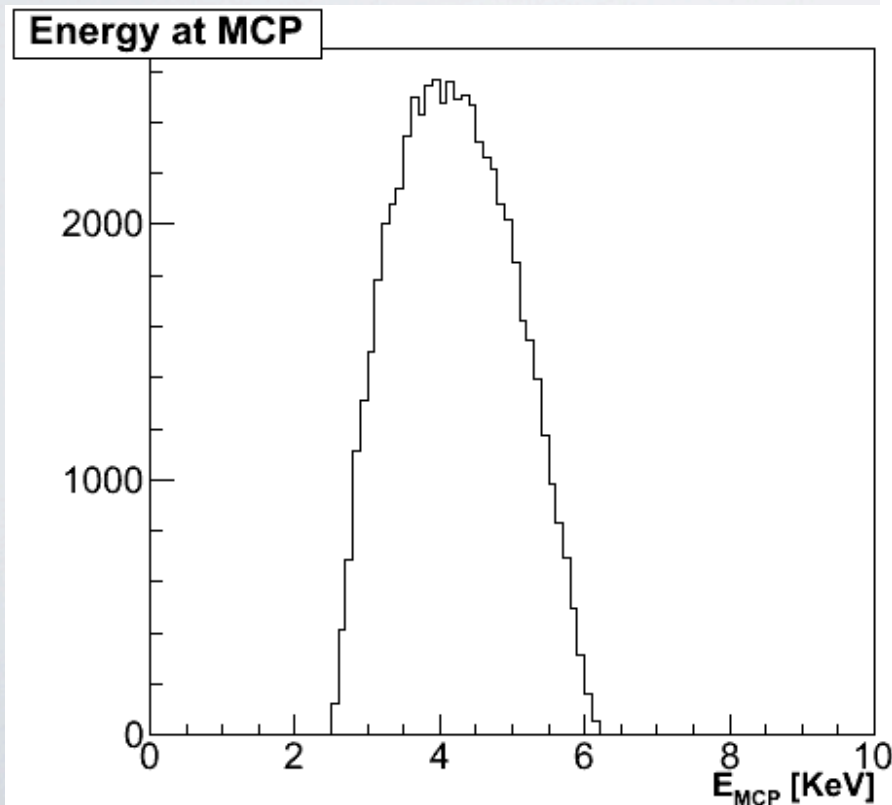
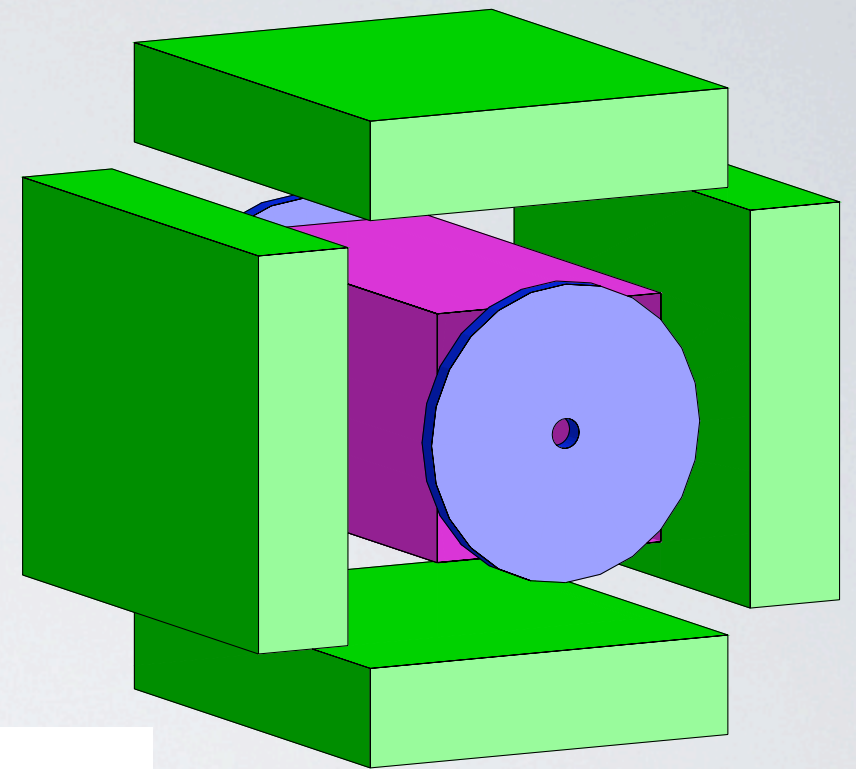


# Some Simulation Results

$\beta$  decay of  $^{21}\text{Na}^+ \rightarrow ^{21}\text{Ne}^+ + \beta^+ + \nu$

Kinetic energy of original ion 4.2 KeV

Cuts require hits in only one MCP and only one SSD (no ambiguity)



# Polarization in $\beta$ -Decay Studies

*Neat trick*

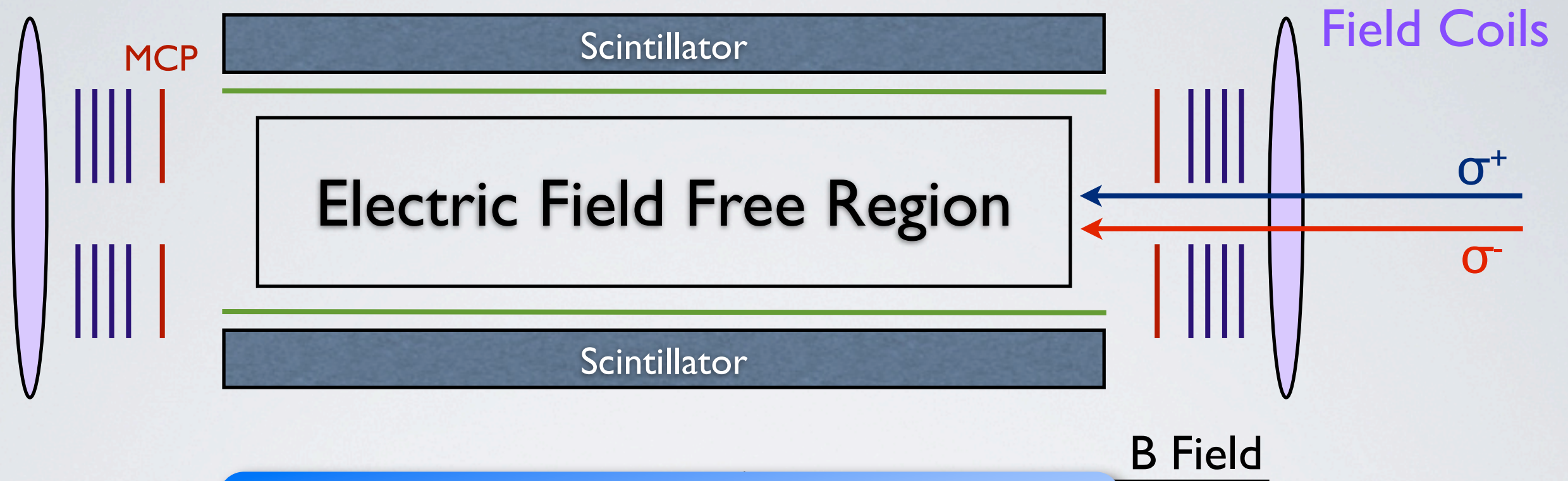


- Add on-axis magnetic field for Zeeman splitting.
- On-axis field does not effect the trajectories ( $\mathbf{V} \times \mathbf{B} = 0$ ).
- Polarize ions with circularly polarized lasers.
- Due to large doppler shift (high energy ions)  $\rightarrow$  two independent ion populations (parallel/ anti-parallel).
- MCP hit is determined by direction of ion  $\rightarrow$  each MCP sees only one population.
- Need polarizable ions (usually singly ionized alkaline earth metals - which look like alkali metals when singly ionized).



# Polarization in $\beta$ -Decay Studies

*Neat trick*



- Add on-axis
- On-axis field
- Polarize ion
- Due to large doppler shift (high energy ions)  $\rightarrow$  two independent ion populations (parallel/anti-parallel).
- MCP hit is determined by direction of ion  $\rightarrow$  each MCP sees only one population.
- Need polarizable ions (usually singly ionized alkaline earth metals - which look like alkali metals when singly ionized).

*Pretty much only possible in  
optical and ES traps*



# EXAMPLE $^{39}\text{K}$

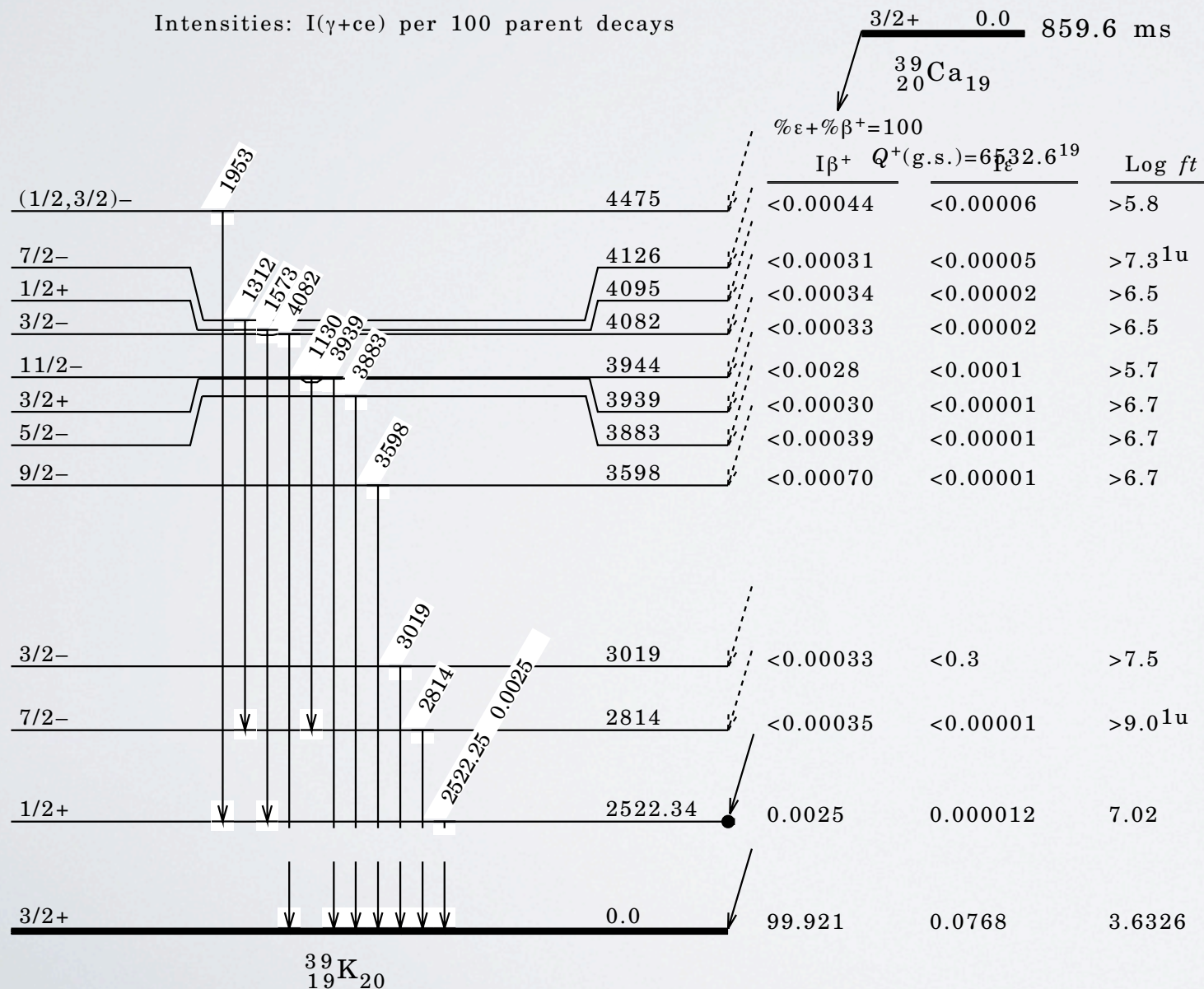
?

$$\frac{d\Gamma}{dE_\beta d\Omega_\beta d\Omega_\nu} \propto \xi \left\{ 1 + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + b \frac{m}{E_e} + c \left[ \frac{1}{3} \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} - \frac{(\vec{p}_e \cdot \vec{j})(\vec{p}_\nu \cdot \vec{j})}{E_e E_\nu} \right] \right. \\ \left. \left[ \frac{J(J+1) - 3 \langle (\vec{J} \cdot \vec{j})^2 \rangle}{J(2J-1)} \right] + \frac{\langle \vec{J} \rangle}{J} \cdot \left[ A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right\}$$

Translates to TOF  
on MCPs

## Decay Scheme

Intensities: I( $\gamma$ +ce) per 100 parent decays



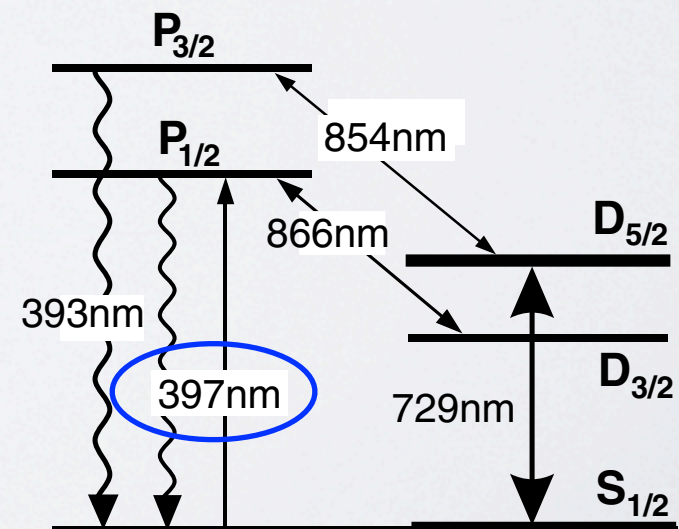
Mirror decay  $3/2^+ \rightarrow 3/2^+$

Combination GF and F but calculable

$Q = 6.532 \text{ MeV}$

$\tau_{1/2} = 860 \text{ ms}$

For 4.2keV ~360GHz doppler shift →  
720GHz  $\sigma_+/\sigma_-$  separation.





# “High Energy” Chemistry (IV)

## *An experiment looking for a theorist*

- Production of radioactive molecules / dimer / clusters is fairly trivial (usually a side effect of the production of radioactive atoms / ions).
- Ionization of such molecules - also trivial.
- ES trap can easily trap molecules of hundreds of amu (also used for bio-molecules).
- Radioactive decay dumps a lot of energy and momentum into the decay products.
- Time scale for decay / emission of shakeoff products is effectively instantaneous.



- Detect molecular decays in ES trap.
- Energy / Momentum sharing between decay products (electronic interaction timescales?).
- Angular correlation in decays (potential?).
- High detection efficiency.
- Mass resolution good enough for selection of different numbers of radioactive atoms in clusters

$$\frac{{}^{23}\text{Na}_4 - {}^{21}\text{Na}{}^{23}\text{Na}_3}{{}^4\text{Na}_{23}} \sim \frac{2}{92}$$
$$\frac{{}^{23}\text{Na}_4 - {}^{21}\text{Na}_2{}^{23}\text{Na}_2}{{}^4\text{Na}_{23}} \sim \frac{4}{92}$$

**Trivial**





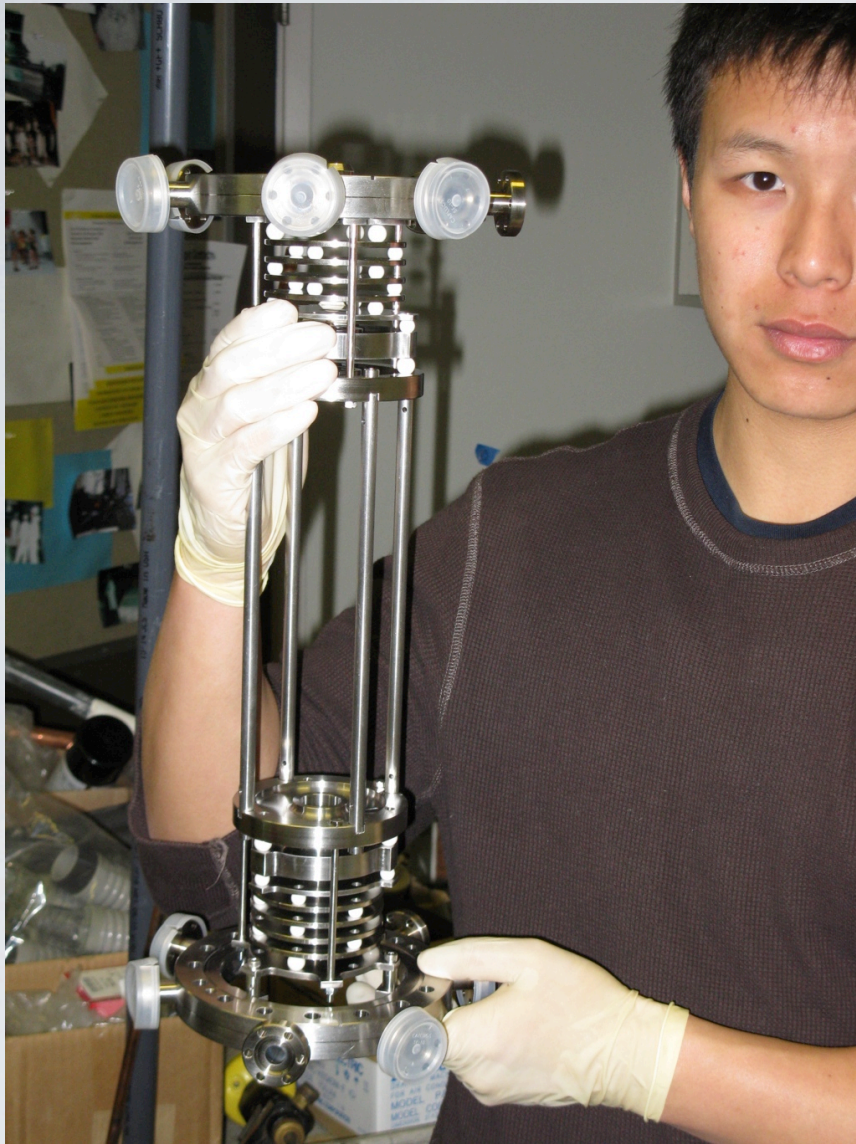
# Open Issues

- Installation at LBL (ECR source, isotope selection (*probably start with  $^{39}\text{Ar}$* )) - looking at using  $^{14}\text{O}$  beamline.
- Detector design: SSD, Thick GEM+Scintillator .
- Detection:
  - Initially we plan to use image charge detection (but that is relevant for large ion populations or highly charged ions).
  - Single ion detection? (Relevant for SH) *SQUID?*
- Ion bunch size:
  - Injection bunch size (fast switches, but would also be nice to bunch first with say an RFQ).
  - WI are trying to reduce bunch size by RF voltages on the electrodes.
  - We want to try several schemes (“stochastic cooling” - RF kicker, Laser cooling where applicable).

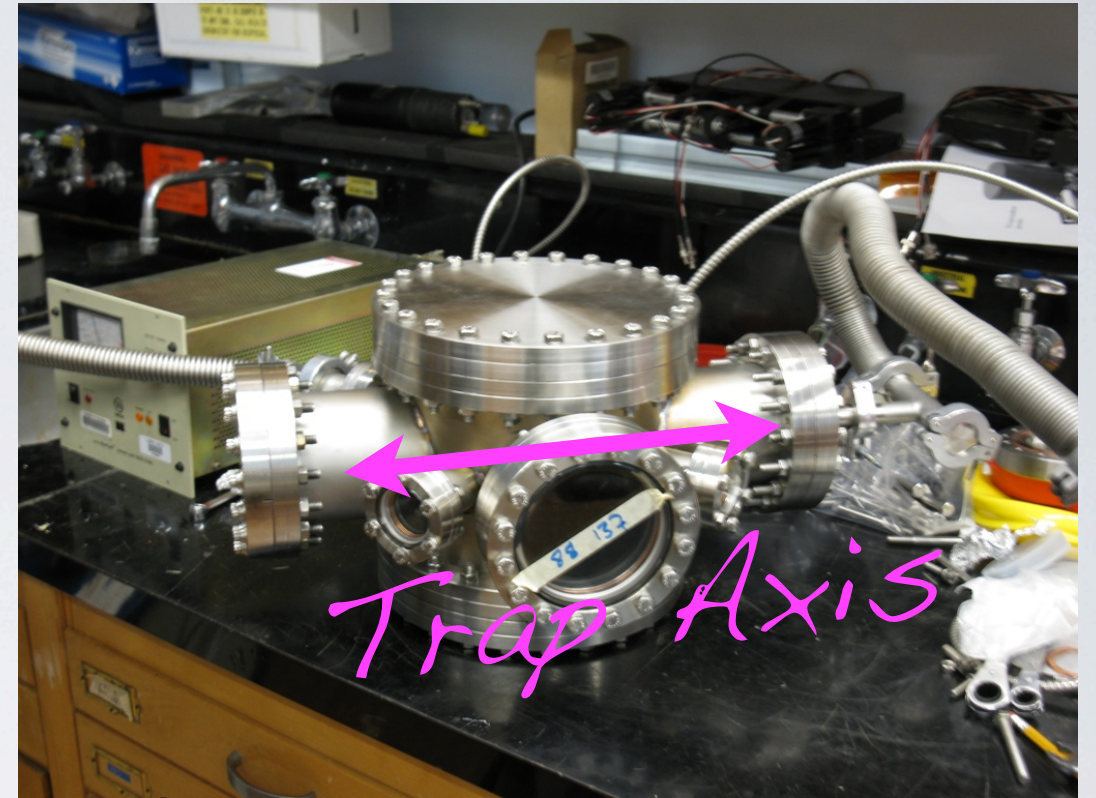




# Status



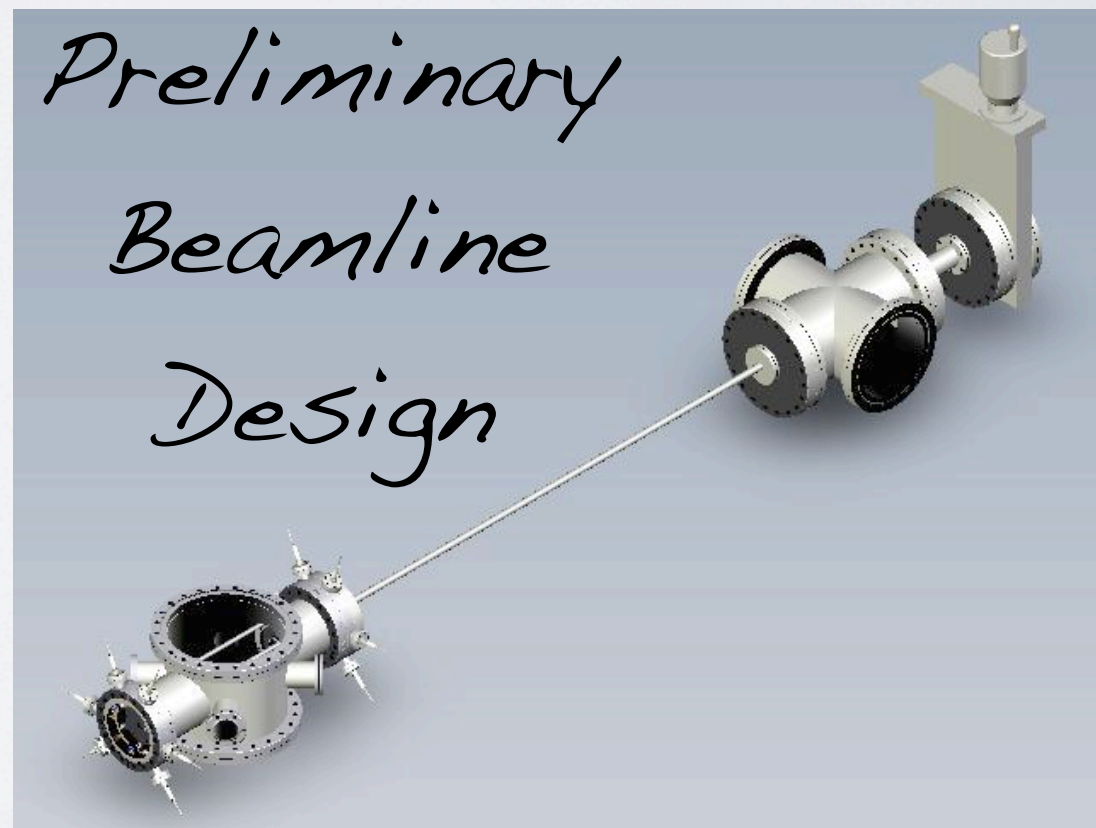
Trap  
during  
fit test



Vacuum Chamber



Trap  
electrode  
sets



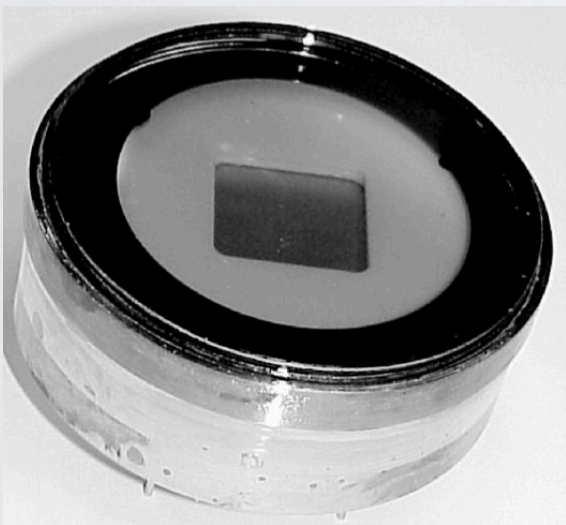
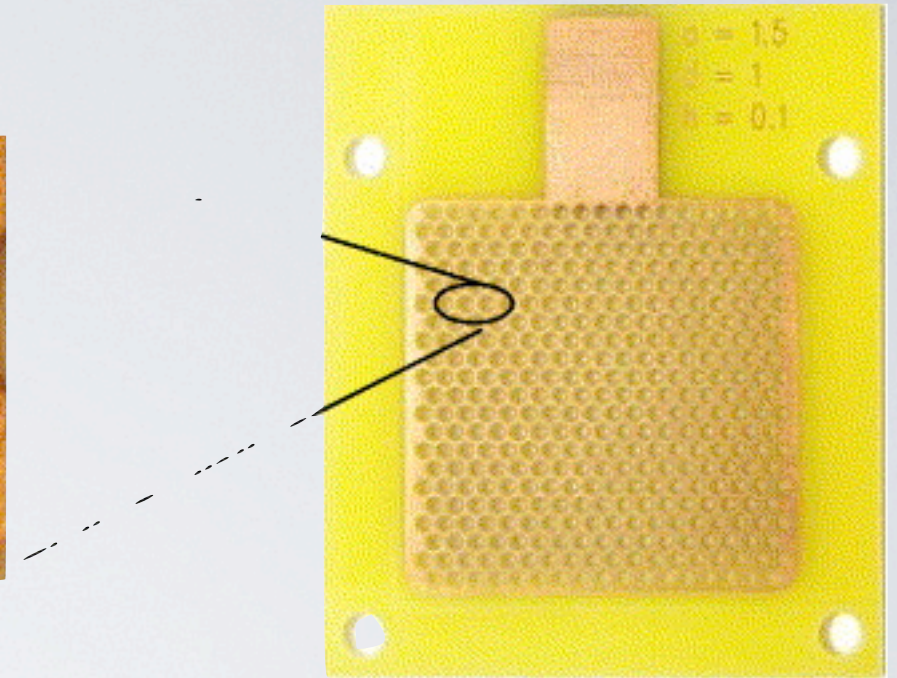
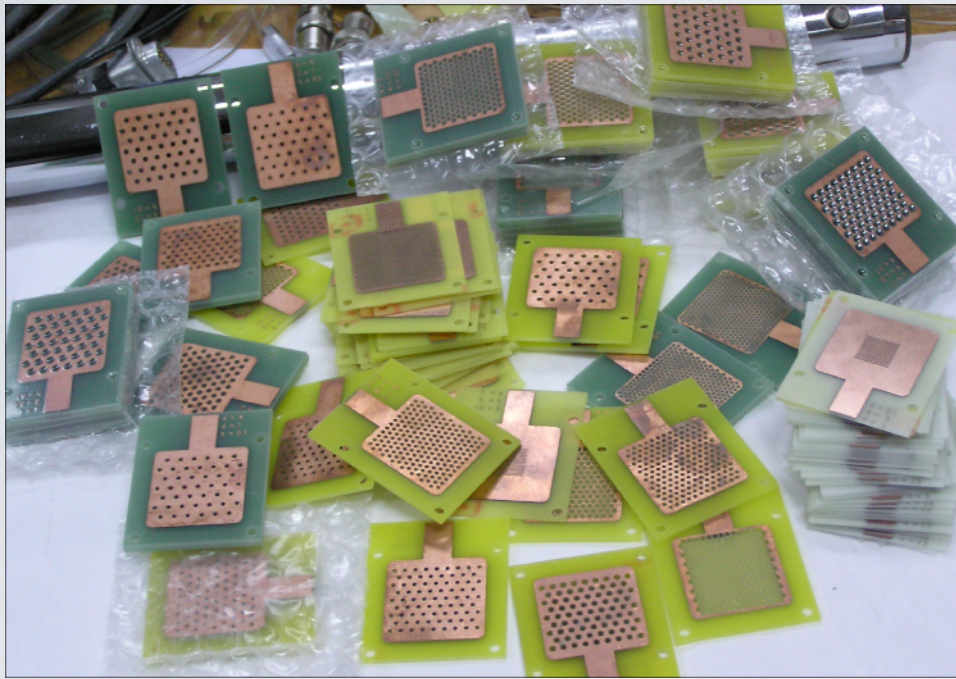


# Summary

- Ion traps are useful for many applications.
- But... conventional ion traps are problematic and technically challenging.
- EIBT is an easy and versatile trap to make. Can almost be a “black box” transportable setup.
- We expect to have a trap setup for installation soon, initially for mass spectroscopy but hopefully later for  $\beta$  decay studies.
- Collaboration / Comments / Ideas always welcome.



# THGEM - Thick Gaseous Electron Multiplier



Sealed GEM (UV Detector)

*M. Balcerzyk et al., IEEE Tran.*

*Nucl. Sci. 50, 1 (2003)*